

A Note on LP -Kenmotsu Manifolds Admitting Conformal Ricci-Yamabe Solitons**Mobin Ahmad^{1,*}, Gazala¹, Maha Atif Al-Shabrawi²**¹*Department of Mathematics and statistics, Integral University, Kursi Road, Lucknow-226026, India*²*Department of Mathematical Sciences, Umm Ul Qura University, Makkah, Saudi Arabia** *Corresponding author: mobinahmad68@gmail.com*

Abstract. In the current note, we study Lorentzian para-Kenmotsu (in brief, LP -Kenmotsu) manifolds admitting conformal Ricci-Yamabe solitons (CRYS) and gradient conformal Ricci-Yamabe soliton (gradient CRYS). At last by constructing a 5-dimensional non-trivial example we illustrate our result.

1. Introduction

As a generalization of the classical Ricci flow [8], the concept of conformal Ricci flow was introduced by Fischer [5], which is defined on an n -dimensional Riemannian manifold M by the equations

$$\frac{\partial g}{\partial t} = -2\left(S + \frac{g}{n}\right) - pg, \quad r(g) = -1,$$

where p defines a time dependent non-dynamical scalar field (also called the conformal pressure), g is the Riemannian metric; r and S represent the scalar curvature and the Ricci tensor of M , respectively. The term $-pg$ plays a role of constraint force to maintain r in the above equation.

In [1], the authors Basu and Bhattacharyya proposed the concept of conformal Ricci soliton on M and is defined by

$$\mathcal{L}_K g + 2S + (2\Lambda - (p + \frac{2}{n}))g = 0,$$

where \mathcal{L}_K represents the Lie derivative operator along the smooth vector field K on M and $\Lambda \in \mathbb{R}$ (the set of real numbers).

Received: Feb. 3, 2023.

2010 *Mathematics Subject Classification*. 53C20, 53C21, 53C25, 53E20.

Key words and phrases. Lorentzian para-Kenmotsu manifolds; conformal Ricci-Yamabe solitons; Einstein manifolds; ν -Einstein manifolds.

Very recently, a scalar combination of Ricci and Yamabe flows was proposed by the authors Güler and Crasmareanu [7], this advanced class of geometric flows called Ricci-Yamabe (RY) flow of type (σ, ρ) and is defined by

$$\frac{\partial}{\partial t}g(t) + 2\sigma S(g(t)) + \rho r(t)g(t) = 0, \quad g(0) = g_0$$

for some scalars σ and ρ . A solution to the RY flow is called Ricci-Yamabe soliton (RYS) if it depends only on one parameter group of diffeomorphism and scaling. A Riemannian (or semi-Riemannian) manifold M is said to have a RYS if [9, 10]

$$\mathcal{L}_K g + 2\sigma S + (2\Lambda - \rho r)g = 0. \quad (1.1)$$

A Riemannian (or semi-Riemannian) manifold M is said to have a conformal Ricci-Yamabe soliton (CRYS) if [20]

$$\mathcal{L}_K g + 2\sigma S + (2\Lambda - \rho r - (p + \frac{2}{n}))g = 0, \quad (1.2)$$

where $\sigma, \rho, \Lambda \in \mathbb{R}$.

If K is the gradient of a smooth function v on M , then (1.2) is called the gradient conformal Ricci-Yamabe soliton (gradient CRYS) and hence (1.2) turns to

$$\nabla^2 v + \sigma S + (\Lambda - \frac{\rho r}{2} - \frac{1}{2}(p + \frac{2}{n}))g = 0, \quad (1.3)$$

where $\nabla^2 v$ is the Hessian of v and is defined by $Hess v = \nabla \nabla v$.

A CRYS is said to be shrinking, steady or expanding if $\Lambda < 0, = 0$ or > 0 , respectively. A CRYS is said to be a

- Conformal Ricci soliton if $\sigma = 1, \rho = 0$,
- Conformal Yamabe soliton if $\sigma = 0, \rho = 1$,
- Conformal Einstein soliton if $\sigma = 1, \rho = -1$.

As a continuation of this study, we tried to study CRYS and gradient CRYS in the frame-work of LP -Kenmotsu manifolds of dimension n . We recommend the papers [2–4, 6, 13–17] and the references therein for more details about the related studies.

2. Preliminaries

An n -dimensional differentiable manifold M with structure (φ, ζ, ν, g) is said to be a Lorentzian almost paracontact metric manifold, if it admits a $(1, 1)$ -tensor field φ , a contravariant vector field ζ , a 1-form ν and a Lorentzian metric g satisfying

$$\nu(\zeta) + 1 = 0, \quad (2.1)$$

$$\varphi^2 E = E + \nu(E)\zeta, \quad (2.2)$$

$$\varphi \zeta = 0, \quad \nu(\varphi E) = 0,$$

$$g(\varphi E, \varphi F) = g(E, F) + \nu(E)\nu(F),$$

$$g(E, \zeta) = \nu(E), \tag{2.3}$$

$$\varphi(E, F) = \varphi(F, E) = g(E, \varphi F)$$

for any vector fields $E, F \in \chi(M)$, where $\chi(M)$ is the Lie algebra of vector fields on M .

If ζ is a killing vector field, the (para) contact structure is called a K -(para) contact. In such a case, we have

$$\nabla_E \zeta = \varphi E.$$

Recently, the authors Haseeb and Prasad defined and studied the following notion:

Definition 2.1. *A Lorentzian almost paracontact manifold M is called Lorentzian para-Kenmostu manifold if [11]*

$$(\nabla_E \varphi)F = -g(\varphi E, F)\zeta - \nu(F)\varphi E$$

for any E, F on M .

In an LP -Kenmostu manifold, we have

$$\nabla_E \zeta = -E - \nu(E)\zeta, \tag{2.4}$$

$$(\nabla_E \nu)F = -g(E, F) - \nu(E)\nu(F), \tag{2.5}$$

where ∇ denotes the Levi-Civita connection respecting to the Lorentzian metric g .

Furthermore, in an LP -Kenmotsu manifold, the following relations hold [11]:

$$g(R(E, F)G, \zeta) = \nu(R(E, F)G) = g(F, G)\nu(E) - g(E, G)\nu(F),$$

$$R(\zeta, E)F = -R(E, \zeta)F = g(E, F)\zeta - \nu(F)E,$$

$$R(E, F)\zeta = \nu(F)E - \nu(E)F,$$

$$R(\zeta, E)\zeta = E + \nu(E)\zeta, \tag{2.6}$$

$$S(E, \zeta) = (n - 1)\nu(E), \quad S(\zeta, \zeta) = -(n - 1), \tag{2.7}$$

$$Q\zeta = (n - 1)\zeta,$$

for any $E, F, G \in \chi(M)$, where R, S and Q represent the curvature tensor, the Ricci tensor and the Q Ricci operator, respectively.

Definition 2.2. [19] *An LP -Kenmotsu manifold M is said to be ν -Einstein manifold if its $S(\neq 0)$ is of the form*

$$S(E, F) = a g(E, F) + b \nu(E)\nu(F),$$

where a and b are smooth functions on M . In particular, if $b = 0$, then M is termed as an Einstein manifold.

Remark 2.1. [12] In an LP-Kenmotsu manifold of n -dimension, S is of the form

$$S(E, F) = \left(\frac{r}{n-1} - 1\right)g(E, F) + \left(\frac{r}{n-1} - n\right)\nu(E)\nu(F), \quad (2.8)$$

where r is the scalar curvature of the manifold.

Lemma 2.1. In an n -dimensional LP-Kenmotsu manifold, we have

$$\zeta(r) = 2(r - n(n-1)), \quad (2.9)$$

$$(\nabla_E Q)\zeta = QE - (n-1)E, \quad (2.10)$$

$$(\nabla_\zeta Q)E = 2QE - 2(n-1)E, \quad (2.11)$$

for any E on M .

Proof. Equation (2.8) yields

$$QE = \left(\frac{r}{n-1} - 1\right)E + \left(\frac{r}{n-1} - n\right)\nu(E)\zeta. \quad (2.12)$$

Taking the covariant derivative of (2.12) with respect to F and making use of (2.4) and (2.5), we lead to

$$(\nabla_F Q)E = \frac{F(r)}{n-1}(E + \nu(E)\zeta) - \left(\frac{r}{n-1} - n\right)(g(E, F)\zeta + \nu(E)F + 2\nu(E)\nu(F)\zeta).$$

By contracting F in the foregoing equation and using trace $\{F \rightarrow (\nabla_F Q)E\} = \frac{1}{2}E(r)$, we find

$$\frac{n-3}{2(n-1)}E(r) = \left\{ \frac{\zeta(r)}{n-1} - (r - n(n-1)) \right\} \nu(E),$$

which by replacing E by ζ and using (2.1) gives (2.9). We refer the readers to see [13] for the proof of (2.10) and (2.11). \square

Remark 2.2. From the equation (2.9), it is noticed that if an n -dimensional LP-Kenmotsu manifold possesses the constant scalar curvature, then $r = n(n-1)$ and hence (2.8) reduces to $S(E, F) = (n-1)g(E, F)$. Thus the manifold under consideration is an Einstein manifold.

3. CRYS on LP-Kenmotsu Manifolds

Let the metric of an n -dimensional LP-Kenmotsu manifold be a conformal Ricci-Yamabe soliton, thus (1.2) holds. By differentiating (1.2) covariantly with respect to G , we have

$$(\nabla_G \mathcal{L}_K g)(E, F) = -2\sigma(\nabla_G S)(E, F) + \rho(Gr)g(E, F). \quad (3.1)$$

Since $\nabla g = 0$, then the following formula [18]

$$(\mathcal{L}_K \nabla_E g - \nabla_E \mathcal{L}_K g - \nabla_{[K, E]} g)(F, G) = -g((\mathcal{L}_K \nabla)(E, F), G) - g((\mathcal{L}_K \nabla)(E, G), F)$$

turns to

$$(\nabla_E \mathcal{L}_K g)(F, G) = g((\mathcal{L}_K \nabla)(E, F), G) + g((\mathcal{L}_K \nabla)(E, G), F).$$

Since the operator $\mathcal{L}_K \nabla$ is symmetric, therefore we have

$$2g((\mathcal{L}_K \nabla)(E, F), G) = (\nabla_E \mathcal{L}_K g)(F, G) + (\nabla_F \mathcal{L}_K g)(E, G) - (\nabla_G \mathcal{L}_K g)(E, F),$$

which by using (3.1) takes the form

$$2g((\mathcal{L}_K \nabla)(E, F), G) = -2\sigma[(\nabla_E S)(F, G) + (\nabla_F S)(G, E) - (\nabla_G S)(E, F)] + \rho[(Er)g(F, G) + (Fr)g(G, E) - (Gr)g(E, F)]. \tag{3.2}$$

Putting $F = \zeta$ in (3.2) and using (2.3), we find

$$2g((\mathcal{L}_K \nabla)(E, \zeta), G) = -2\sigma[(\nabla_E S)(\zeta, G) + (\nabla_\zeta S)(G, E) - (\nabla_G S)(E, \zeta)] + \rho[(Er)\nu(G) + 2(r - n(n - 1))g(E, G) - (Gr)\nu(E)]. \tag{3.3}$$

By virtue of (2.10) and (2.11), (3.3) leads to

$$2g((\mathcal{L}_K \nabla)(E, \zeta), G) = -4\sigma[S(E, G) - (n - 1)g(E, G)] + \rho[(Er)\nu(G) + 2(r - n(n - 1))g(E, G) - (Gr)\nu(E)].$$

By eliminating G from the foregoing equation, we have

$$2(\mathcal{L}_K \nabla)(F, \zeta) = \rho g(Dr, F)\zeta - \rho(Dr)\nu(F) - 4\sigma QF + [4\sigma(n - 1) + 2\rho(r - n(n - 1))]F. \tag{3.4}$$

If we take r as constant, then from (2.9) it follows that $r = n(n - 1)$, and hence (3.4) reduces to

$$(\mathcal{L}_K \nabla)(F, \zeta) = -2\sigma QF + 2\sigma(n - 1)F. \tag{3.5}$$

Taking covariant derivative of (3.5) with respect to E , we have

$$(\nabla_E \mathcal{L}_K \nabla)(F, \zeta) = (\mathcal{L}_K \nabla)(F, E) - 2\sigma\nu(E)[QF - (n - 1)F] - 2\sigma(\nabla_E Q)F. \tag{3.6}$$

Again from [18], we have

$$(\mathcal{L}_K R)(E, F)G = (\nabla_E \mathcal{L}_K \nabla)(F, G) - (\nabla_F \mathcal{L}_K \nabla)(E, G),$$

which by putting $G = \zeta$ and using (3.6) takes the form

$$(\mathcal{L}_K R)(E, F)\zeta = 2\sigma\nu(F)(QE - (n - 1)E) - 2\sigma\nu(E)(QF - (n - 1)F) - 2\sigma((\nabla_E Q)F - (\nabla_F Q)E). \tag{3.7}$$

Putting $F = \zeta$ in (3.7) then using (2.1), (2.2), (2.10) and (2.11), we arrive at

$$(\mathcal{L}_K R)(E, \zeta)\zeta = 0. \tag{3.8}$$

The Lie derivative of (2.6) along K leads to

$$(\mathcal{L}_K R)(E, \zeta)\zeta - g(E, \mathcal{L}_K \zeta)\zeta + 2\nu(\mathcal{L}_K \zeta)E = -(\mathcal{L}_K \nu)(E)\zeta. \tag{3.9}$$

From (3.8) and (3.9), we have

$$(\mathcal{L}_K \nu)(E)\zeta = -2\nu(\mathcal{L}_K \zeta)E + g(E, \mathcal{L}_K \zeta)\zeta. \quad (3.10)$$

Taking the Lie derivative of $g(E, \zeta) = \nu(E)$, we find

$$(\mathcal{L}_K \nu)(E) = g(E, \mathcal{L}_K \zeta) + (\mathcal{L}_K g)(E, \zeta). \quad (3.11)$$

By putting $F = \zeta$ in (1.2) and using (2.7), we have

$$(\mathcal{L}_K g)(E, \zeta) = -\{2\sigma(n-1) + 2\Lambda - \rho n(n-1) - (p + \frac{2}{n})\}\nu(E), \quad (3.12)$$

where $r = n(n-1)$ being used.

Taking the Lie derivative of $g(\zeta, \zeta) = -1$ along K we lead to

$$(\mathcal{L}_K g)(\zeta, \zeta) = -2\nu(\mathcal{L}_K \zeta). \quad (3.13)$$

From (3.12) and (3.13), we find

$$\nu(\mathcal{L}_K \zeta) = -\{\sigma(n-1) + \Lambda - \frac{\rho n(n-1)}{2} - \frac{1}{2}(p + \frac{2}{n})\}. \quad (3.14)$$

Now combining the equations (3.10), (3.11), (3.12) and (3.14), we find

$$\Lambda = \frac{\rho n(n-1)}{2} - \sigma(n-1) + \frac{1}{2}(p + \frac{2}{n}). \quad (3.15)$$

Thus we have

Theorem 3.1. *Let (M, g) be an n -dimensional LP-Kenmotsu manifold admitting CRYs with constant scalar curvature tensor, then $\Lambda = \frac{\rho n(n-1)}{2} - \sigma(n-1) + \frac{1}{2}(p + \frac{2}{n})$.*

Corollary 3.1. *Let the metric of n -dimensional LP-Kenmotsu manifold is CRYs. Then we have*

Values of σ, ρ	Soliton type	Soliton constant	CRYs to be expanding, shrinking or steady
$\sigma = 1, \rho = 0$	conformal Ricci soliton	$\Lambda = \frac{1}{2}(p + \frac{2}{n}) - (n-1)$	CRYs is shrinking, steady and expanding if $p > \frac{2(n^2-n-1)}{n}$, $p = \frac{2(n^2-n-1)}{n}$ and $p < \frac{2(n^2-n-1)}{n}$, resp.
$\sigma = 0, \rho = 1$	conformal Yamabe soliton	$\Lambda = \frac{1}{2}(p + \frac{2}{n}) + \frac{n(n-1)}{2}$	CRYs is shrinking, steady and expanding if $p < \frac{-(n^3-n^2+2)}{n}$, $p = \frac{-(n^3-n^2+2)}{n}$ and $p > \frac{-(n^3-n^2+2)}{n}$, resp.
$\sigma = 1, \rho = -1$	conformal Einstein soliton	$\Lambda = \frac{1}{2}(p + \frac{2}{n}) - \frac{(n-1)(n+2)}{2}$	CRYs is shrinking, steady and expanding if $p < \frac{(n+1)(n^2-2)}{n}$, $p = \frac{(n+1)(n^2-2)}{n}$ and $p > \frac{(n+1)(n^2-2)}{n}$, resp.

4. Gradient CRYS on LP -Kenmotsu Manifolds

Let M be an n -dimensional LP -Kenmotsu manifold with g as a gradient CRYS. Then equation (1.3) can be written as

$$\nabla_E Dv + \sigma QE + (\Lambda - \frac{\rho r}{2} - \frac{1}{2}(\rho + \frac{2}{n}))E = 0, \tag{4.1}$$

for all vector fields E on M , where D denotes the gradient operator of g . Taking the covariant derivative of (4.1) with respect to F , we have

$$\begin{aligned} \nabla_F \nabla_E Dv &= -\sigma\{(\nabla_F Q)E + Q(\nabla_F E)\} + \rho \frac{F(r)}{2} E \\ &\quad - (\Lambda - \frac{\rho r}{2} - \frac{1}{2}(\rho + \frac{2}{n}))\nabla_F E. \end{aligned} \tag{4.2}$$

Interchanging E and F in (4.2), we lead to

$$\begin{aligned} \nabla_E \nabla_F Dv &= -\sigma\{(\nabla_E Q)F + Q(\nabla_E F)\} + \rho \frac{E(r)}{2} F \\ &\quad - (\Lambda - \frac{\rho r}{2} - \frac{1}{2}(\rho + \frac{2}{n}))\nabla_E F. \end{aligned} \tag{4.3}$$

By making use of (4.1)-(4.3), we find

$$R(E, F)Dv = \sigma\{(\nabla_F Q)E - (\nabla_E Q)F\} + \frac{\rho}{2}\{E(r)F - F(r)E\}. \tag{4.4}$$

Now from (2.8), we find

$$QE = (\frac{r}{n-1} - 1)E + (\frac{r}{n-1} - n)\nu(E)\zeta,$$

which on taking covariant derivative with respect to F leads to

$$\begin{aligned} (\nabla_F Q)E &= \frac{F(r)}{n-1}(E + \nu(E)\zeta) - (\frac{r}{n-1} - n)(g(E, F)\zeta \\ &\quad + 2\nu(E)\nu(F)\zeta + \nu(E)F). \end{aligned} \tag{4.5}$$

By using (4.5) in (4.4), we have

$$\begin{aligned} R(E, F)Dv &= \frac{(n-1)\rho - 2\sigma}{2(n-1)}\{E(r)F - F(r)E\} + \frac{\sigma}{n-1}\{F(r)\nu(E)\zeta - E(r)\nu(F)\zeta\} \\ &\quad - \sigma(\frac{r}{n-1} - n)(\nu(E)F - \nu(F)E). \end{aligned} \tag{4.6}$$

Contracting forgoing equation along E gives

$$\begin{aligned} S(F, Dv) &= -\left\{\frac{(n-1)^2\rho - 2\sigma(n-2)}{2(n-1)}\right\}F(r) \\ &\quad + \frac{\sigma(n-3)(r - n(n-1))}{n-1}\nu(F). \end{aligned} \tag{4.7}$$

From the equation (2.8), we have

$$S(F, Dv) = (\frac{r}{n-1} - 1)F(v) + (\frac{r}{n-1} - n)\nu(F)\zeta(v). \tag{4.8}$$

Now by equating (4.7) and (4.8), then putting $F = \zeta$ and using (2.1), (2.9), we find

$$\zeta(v) = \frac{r - n(n-1)}{n-1} \{ \sigma - (n-1)\rho \}. \quad (4.9)$$

Taking the inner product of (4.6) with ζ , we get

$$F(v)\nu(E) - E(v)\nu(F) = \frac{\rho}{2} \{ E(r)\nu(F) - F(r)\nu(E) \},$$

which by replacing E by ζ then using (2.9) and (4.9), we infer

$$F(v) = -\frac{\sigma(r - n(n-1))}{n-1} \nu(F) - \frac{\rho}{2} F(r). \quad (4.10)$$

If we take r as constant, then from Remark 2.5, we get $r = n(n-1)$. Thus (4.10) leads to $F(v) = 0$. This implies that v is constant. Thus the soliton under the consideration is trivial. Hence we state:

Theorem 4.1. *If the metric of an n -dimensional LP-Kenmotsu manifold of constant scalar curvature tensor admitting a special type of vector field is gradient CRYS, then the soliton is trivial.*

For v constant, (1.3) turns to

$$\sigma QE = -\left(\Lambda - \frac{\rho r}{2} - \frac{1}{2} \left(p + \frac{2}{n} \right) \right) E,$$

which leads to

$$S(E, F) = -\frac{1}{\sigma} \left(\Lambda - \frac{\rho n(n-1)}{2} - \frac{1}{2} \left(p + \frac{2}{n} \right) \right) g(E, F), \quad \sigma \neq 0. \quad (4.11)$$

By putting $E = F = \zeta$ in (4.11) and using (2.7), we obtain

$$\Lambda = \frac{\rho n(n-1)}{2} - \sigma(n-1) + \frac{1}{2} \left(p + \frac{2}{n} \right). \quad (4.12)$$

Corollary 4.1. *If an n -dimensional LP-Kenmotsu manifold admits a gradient CRYS with the constant scalar curvature, then the manifold under the consideration is an Einstein manifold and $\Lambda = \frac{\rho n(n-1)}{2} - \sigma(n-1) + \frac{1}{2} \left(p + \frac{2}{n} \right)$.*

5. Example

We consider the 5-dimensional manifold $M^5 = \{ (x_1, x_2, x_3, x_4, x_5) \in \mathbb{R}^5 : x_5 > 0 \}$, where $(x_1, x_2, x_3, x_4, x_5)$ are the standard coordinates in \mathbb{R}^5 . Let $\varrho_1, \varrho_2, \varrho_3, \varrho_4$ and ϱ_5 be the vector fields on M^5 given by

$$\varrho_1 = e^{x_5} \frac{\partial}{\partial x_1}, \quad \varrho_2 = e^{x_5} \frac{\partial}{\partial x_2}, \quad \varrho_3 = e^{x_5} \frac{\partial}{\partial x_3}, \quad \varrho_4 = e^{x_5} \frac{\partial}{\partial x_4}, \quad \varrho_5 = \frac{\partial}{\partial x_5} = \zeta,$$

which are linearly independent at each point of M^5 . Let g be the Lorentzian metric defined by

$$g(\varrho_i, \varrho_i) = 1, \quad \text{for } 1 \leq i \leq 4 \quad \text{and} \quad g(\varrho_5, \varrho_5) = -1,$$

$$g(\varrho_i, \varrho_j) = 0, \quad \text{for } i \neq j, \quad 1 \leq i, j \leq 5.$$

Let ν be the 1-form defined by $\nu(E) = g(E, \varrho_5) = g(E, \zeta)$ for all $E \in \chi(M^5)$, and let φ be the $(1, 1)$ -tensor field defined by

$$\varphi\varrho_1 = -\varrho_2, \varphi\varrho_2 = -\varrho_1, \varphi\varrho_3 = -\varrho_4, \varphi\varrho_4 = -\varrho_3, \varphi\varrho_5 = 0.$$

By applying linearity of φ and g , we have

$$\nu(\zeta) = g(\zeta, \zeta) = -1, \varphi^2 E = E + \nu(E)\zeta \text{ and } g(\varphi E, \varphi F) = g(E, F) + \nu(E)\nu(F)$$

for all $E, F \in \chi(M^5)$. Thus for $\varrho_5 = \zeta$, the structure (φ, ζ, ν, g) defines a Lorentzian almost paracontact metric structure on M^5 . Then we have

$$[\varrho_i, \varrho_j] = -\varrho_i, \quad \text{for } 1 \leq i \leq 4, j = 5,$$

$$[\varrho_i, \varrho_j] = 0, \quad \text{otherwise.}$$

By using Koszul's formula, we can easily find we obtain

$$\nabla_{\varrho_i} \varrho_j = \begin{cases} -\varrho_5, & 1 \leq i = j \leq 4, \\ -\varrho_i, & 1 \leq i \leq 4, j = 5, \\ 0, & \text{otherwise.} \end{cases}$$

Also one can easily verify that

$$\nabla_E \zeta = -E - \eta(E)\zeta \quad \text{and} \quad (\nabla_E \varphi)F = -g(\varphi E, F)\zeta - \nu(F)\varphi E.$$

Therefore, the manifold is an LP -Kenmotsu manifold.

From the above results, we can easily obtain the non-vanishing components of R as follows:

$$\begin{aligned} R(\varrho_1, \varrho_2)\varrho_1 &= -\varrho_2, \quad R(\varrho_1, \varrho_2)\varrho_2 = \varrho_1, \quad R(\varrho_1, \varrho_3)\varrho_1 = -\varrho_3, \quad R(\varrho_1, \varrho_3)\varrho_3 = \varrho_1, \\ R(\varrho_1, \varrho_4)\varrho_1 &= -\varrho_4, \quad R(\varrho_1, \varrho_4)\varrho_4 = \varrho_1, \quad R(\varrho_1, \varrho_5)\varrho_1 = -\varrho_5, \quad R(\varrho_1, \varrho_5)\varrho_5 = -\varrho_1, \\ R(\varrho_2, \varrho_3)\varrho_2 &= -\varrho_3, \quad R(\varrho_2, \varrho_3)\varrho_3 = \varrho_2, \quad R(\varrho_2, \varrho_4)\varrho_2 = -\varrho_4, \quad R(\varrho_2, \varrho_4)\varrho_4 = \varrho_2, \\ R(\varrho_2, \varrho_5)\varrho_2 &= -\varrho_5, \quad R(\varrho_2, \varrho_5)\varrho_5 = -\varrho_2, \quad R(\varrho_3, \varrho_4)\varrho_3 = -\varrho_4, \quad R(\varrho_3, \varrho_4)\varrho_4 = \varrho_3, \\ R(\varrho_3, \varrho_5)\varrho_3 &= -\varrho_5, \quad R(\varrho_3, \varrho_5)\varrho_5 = -\varrho_3, \quad R(\varrho_4, \varrho_5)\varrho_4 = -\varrho_5, \quad R(\varrho_4, \varrho_5)\varrho_5 = -\varrho_4. \end{aligned}$$

Also, we calculate the Ricci tensors as follows:

$$S(\varrho_1, \varrho_1) = S(\varrho_2, \varrho_2) = S(\varrho_3, \varrho_3) = S(\varrho_4, \varrho_4) = 4, \quad S(\varrho_5, \varrho_5) = -4.$$

Therefore, we have

$$r = S(\varrho_1, \varrho_1) + S(\varrho_2, \varrho_2) + S(\varrho_3, \varrho_3) + S(\varrho_4, \varrho_4) - S(\varrho_5, \varrho_5) = 20.$$

Now by taking $D\nu = (\varrho_1\nu)\varrho_1 + (\varrho_2\nu)\varrho_2 + (\varrho_3\nu)\varrho_3 + (\varrho_4\nu)\varrho_4 + (\varrho_5\nu)\varrho_5$, we have

$$\begin{aligned} \nabla_{\varrho_1} D\nu &= (\varrho_1(\varrho_1\nu) - (\varrho_5\nu))\varrho_1 + (\varrho_1(\varrho_2\nu))\varrho_2 + (\varrho_1(\varrho_3\nu))\varrho_3 + (\varrho_1(\varrho_4\nu))\varrho_4 \\ &\quad + (\varrho_1(\varrho_5\nu) - (\varrho_1\nu))\varrho_5, \end{aligned}$$

$$\begin{aligned}\nabla_{\varrho_2} Dv &= (\varrho_2(\varrho_1 v))\varrho_1 + (\varrho_2(\varrho_2 v) - (\varrho_5 v))\varrho_2 + (\varrho_2(\varrho_3 v))\varrho_3 + (\varrho_2(\varrho_4 v))\varrho_4 \\ &\quad + (\varrho_2(\varrho_5 v) - (\varrho_2 v))\varrho_5,\end{aligned}$$

$$\begin{aligned}\nabla_{\varrho_3} Dv &= (\varrho_3(\varrho_1 v))\varrho_1 + (\varrho_3(\varrho_2 v))\varrho_2 + (\varrho_3(\varrho_3 v) - (\varrho_5 v))\varrho_3 + (\varrho_3(\varrho_4 v))\varrho_4 \\ &\quad + (\varrho_3(\varrho_5 v) - (\varrho_3 v))\varrho_5,\end{aligned}$$

$$\begin{aligned}\nabla_{\varrho_4} Dv &= (\varrho_4(\varrho_1 v))\varrho_1 + (\varrho_4(\varrho_2 v))\varrho_2 + (\varrho_4(\varrho_3 v))\varrho_3 + (\varrho_4(\varrho_4 v) - (\varrho_5 v))\varrho_4 \\ &\quad + (\varrho_4(\varrho_5 v) - (\varrho_4 v))\varrho_5,\end{aligned}$$

$$\nabla_{\varrho_5} Dv = (\varrho_5(\varrho_1 v))\varrho_1 + (\varrho_5(\varrho_2 v))\varrho_2 + (\varrho_5(\varrho_3 v))\varrho_3 + (\varrho_5(\varrho_4 v))\varrho_4 + (\varrho_5(\varrho_5 v))\varrho_5.$$

Thus by virtue of (4.1), we obtain

$$\left\{ \begin{array}{l} \varrho_1(\varrho_1 v) - \varrho_5 v = -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \varrho_2(\varrho_2 v) - \varrho_5 v = -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \varrho_3(\varrho_3 v) - \varrho_5 v = -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \varrho_4(\varrho_4 v) - \varrho_5 v = -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \varrho_5(\varrho_5 v) = -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \varrho_1(\varrho_2 v) = \varrho_1(\varrho_3 v) = \varrho_1(\varrho_4 v) = 0, \\ \varrho_2(\varrho_1 v) = \varrho_2(\varrho_3 v) = \varrho_2(\varrho_4 v) = 0, \\ \varrho_3(\varrho_1 v) = \varrho_3(\varrho_2 v) = \varrho_3(\varrho_4 v) = 0, \\ \varrho_4(\varrho_1 v) = \varrho_4(\varrho_2 v) = \varrho_4(\varrho_3 v) = 0, \\ \varrho_1(\varrho_5 v) - (\varrho_1 v) = \varrho_2(\varrho_5 v) - (\varrho_2 v) = 0, \\ \varrho_3(\varrho_5 v) - (\varrho_3 v) = \varrho_4(\varrho_5 v) - (\varrho_4 v) = 0. \end{array} \right. \quad (5.1)$$

Thus the equations in (5.1) are respectively amounting to

$$\begin{aligned}e^{2x_5} \frac{\partial^2 v}{\partial x_1^2} - \frac{\partial v}{\partial x_5} &= -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ e^{2x_5} \frac{\partial^2 v}{\partial x_2^2} - \frac{\partial v}{\partial x_5} &= -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ e^{2x_5} \frac{\partial^2 v}{\partial x_3^2} - \frac{\partial v}{\partial x_5} &= -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ e^{2x_5} \frac{\partial^2 v}{\partial x_4^2} - \frac{\partial v}{\partial x_5} &= -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \frac{\partial^2 v}{\partial x_5^2} &= -(\Lambda + 4\sigma - 10\rho - \frac{1}{2}(p + \frac{2}{5})), \\ \frac{\partial^2 v}{\partial x_1 \partial x_2} = \frac{\partial^2 v}{\partial x_1 \partial x_3} = \frac{\partial^2 v}{\partial x_1 \partial x_4} = \frac{\partial^2 v}{\partial x_2 \partial x_3} = \frac{\partial^2 v}{\partial x_2 \partial x_4} = \frac{\partial^2 v}{\partial x_3 \partial x_4} &= 0,\end{aligned}$$

$$e^{x_5} \frac{\partial^2 v}{\partial x_5 \partial x_1} + \frac{\partial v}{\partial x_1} = e^{x_5} \frac{\partial^2 v}{\partial x_5 \partial x_2} + \frac{\partial v}{\partial x_2} = e^{x_5} \frac{\partial^2 v}{\partial x_5 \partial x_3} + \frac{\partial v}{\partial x_3} = e^{x_5} \frac{\partial^2 v}{\partial x_5 \partial x_4} + \frac{\partial v}{\partial x_4} = 0.$$

From the above equations it is observed that v is constant for $\Lambda = -4\sigma + 10\rho + \frac{1}{2}(p + \frac{2}{5})$. Hence equation (4.1) is satisfied. Thus, g is a gradient RYS with the soliton vector field $K = Dv$, where v is constant and $\Lambda = -4\sigma + 10\rho + \frac{1}{2}(p + \frac{2}{5})$. Thus, Theorem 4.1 is verified.

Conflicts of Interest: The authors declare that there are no conflicts of interest regarding the publication of this paper.

Acknowledgement: The first and second authors would like to thank the Integral University, Lucknow, India, for providing the manuscript number IU/R&D/2022-MCN0001737 to the present work.

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