International Journal of Analysis and Applications

Complex Dynamics of a Predator-Prey System With Gompertz Growth and Herd Behavior

Rizwan Ahmed¹, M. B. Almatrafi^{2,*}

¹Department of Mathematics, Air University Multan Campus, Multan, Pakistan ²Department of Mathematics, College of Science, Taibah University, Al-Madinah Al -Munawarah, Saudi Arabia

* Corresponding author: mmutrafi@taibahu.edu.sa

Abstract. The complex dynamics of a predator-prey system in discrete time are studied. In this system, we consider the prey's Gompertz growth and the square-root functional response. The existence of fixed points and stability are examined. Using the center manifold and bifurcation theory, we found that the system undergoes transcritical bifurcation, period-doubling bifurcation, and Neimark-Sacker bifurcation. In addition, numerical examples are presented to illustrate the consistency of the analytical findings. The bifurcation diagrams show that the positive fixed point is stable if the death rate of the predator is greater than a threshold value. Biologically, it means that to prevent the predator population from growing uncontrollably and stability of the positive fixed point, the predator's death rate should be greater than the threshold value.

1. Introduction

In mathematical ecology, a predator-prey interaction is essential due to its significance and universal existence. One of the foremost imperative subjects in mathematical ecology is the dynamic mutual action between predator and prey populations, which aids in conserving species in a habitat. The first predator-prey interaction was derived independently by Lotka [1], and Volterra [2], and it is known as the Lotka-Volterra predator-prey system. Since then, numerous researchers have made much progress considering different biological facts.

Received: Jul. 12, 2023.

²⁰²⁰ Mathematics Subject Classification. 39A28, 39A30, 92D25.

Key words and phrases. predator-prey; functional response; herd behavior; stability; bifurcation.

The functional response is an essential component of predator-prey interactions within population dynamics. It illustrates the link between the predator's consumption rate and the density of prey. It denotes how much prey each predator consumes. Holling (1965) proposed three distinct types of functional responses [3]. Subsequently, a number of functional responses were proposed by scholars such as Crowley-Martin [4] and Beddington-DeAngelis [5, 6]. Later, several research studies investigated models centered on predator-prey interactions, encompassing diverse categories of functional responses [7–12].

Some prey populations exhibit herd behavior, where interactions between predators and prey occur primarily at the boundaries of the prey population. The nature of this interaction is not sufficiently explicable through Holling-type responses. It is noteworthy that a particular population of prey exhibits collective behavior known as herd behavior. Therefore, the rate at which a predator captures its prey differs from conventional models. As an example, the predation rate of zooplankton by fish in marine ecosystems exceeds the predation rate of phytoplankton. The observed phenomenon entails the manifestation of herding behavior among phytoplankton. The utilization of the square root of the prey population was employed by Ajraldi et al. [13] as a method to investigate the herd behavior of the prey population. This strategy facilitated the predator's ability to interact with the prey in the boundary region of the group. Numerous scholars have investigated the mechanisms of predator-prey relationships by utilizing square-root functional responses [14–23].

We examine a predator-prey interaction when the prey exhibits group defense, which is represented by the set of ordinary differential equations given as:

$$\begin{cases} \frac{dx}{dt} = xg(x, \mathcal{K}) - yp(x), \\ \frac{dy}{dt} = -dy + q(x)y, \end{cases}$$
(1.1)

where x, y, K > 0, and d > 0 are the prey population density, predator population density, carrying capacity, and predator's death rate, respectively. In the absence of a predator, the function g(x, K) indicates the particular growth rate of the prey. We employ Gompertz growth [24] of prey $g(x, K) = r \ln(\frac{K}{x})$, with a natural growth rate r > 0. The predator response function is represented by p(x), and we assume it is of square root type, i.e, $p(x) = m\sqrt{x}$, where m > 0 represents the predator's search efficiency. Moreover, q(x) describes the conversion rate of prey. For some positive constant c, q(x) = cp(x) in Gauss' model. Considering all assumptions, system (1.1) takes the following form:

$$\begin{cases} \frac{dx}{dt} = rx \ln(\frac{K}{x}) - m\sqrt{x}y, \\ \frac{dy}{dt} = -dy + cm\sqrt{x}y. \end{cases}$$
(1.2)

Due to the square root term, the Jacobian of system (1.2) possesses a singularity. We utilize the transformations $x(t) = u^2(t)$ and y(t) = v(t) to understand it better. After applying this transformation to system (1.2), we obtain

$$\begin{cases} \frac{du}{dt} = \frac{1}{2}ru\ln(\frac{\kappa}{u^2}) - \frac{1}{2}mv, \\ \frac{dv}{dt} = -dv + cmuv. \end{cases}$$
(1.3)

In cases where a population has no overlapping generations, discrete systems that are governed by difference equations are comparatively more appropriate than continuous systems. Several species, such as monocarpic plants and semelparous animals, exhibit distinct generations that remain different from one another, and reproductive events occur at anticipated intervals during specific mating seasons. Discrete-time mappings or difference equations are employed to depict their interrelationships. Discrete-time dynamical systems have complicated and diversified dynamical properties [25–31]. The system to be analyzed in this work is then obtained by using the forward Euler technique on the system (1.3) as follows:

$$\begin{cases} u_{n+1} = u_n + \frac{h}{2} (r u_n \ln(\frac{K}{u_n^2}) - m v_n), \\ v_{n+1} = v_n + h (c m u_n v_n - d v_n), \end{cases}$$
(1.4)

where h > 0 denotes the step size. The main contributions of this study are listed as follows:

- A discrete-time predator-prey system with Gompertz growth and herd behavior is proposed.
- The study investigates the existence and topological categorization of fixed points.
- The results of our study indicate that the system (1.4) undergoes transcritical, period-doubling (PD), and Neimark-Sacker (NS) bifurcations.

The following outlines the format of the paper: Section 2 investigates the existence and stability of fixed points in the system (1.4). In Section 3, we employ the center manifold theorem and bifurcation theory to analyze local bifurcation analysis at fixed points of the system (1.4). Section 4 presents some examples to verify our theoretical results. Some closing observations are added in Section 5.

2. Existence and stability of fixed points

The long-term behavior of dynamical system depends on fixed points. Stable fixed points cause the system to converge, while unstable ones cause oscillation or chaos. The existence and stability conditions for the fixed points of system (1.4) are investigated in this section. By simple algebraic computations, the system (1.4) is found to have two fixed points: $P_1(\sqrt{K}, 0)$ and $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$. The first fixed point P_1 is a boundary point. Biologically, it indicates that when there are no predators around, the number of prey approaches the square root of the carrying capacity. The second fixed point, P_2 is the unique positive fixed point in the system (1.4) if $K > \frac{d^2}{c^2m^2}$.

Now, we will look at the local stability of the fixed points in the system (1.4). The eigenvalues of the Jacobian matrix computed at the fixed points define the local stability of the fixed points. At the point (\bar{u}, \bar{v}) , the Jacobian matrix of system (1.4) is

$$J(\bar{u},\bar{v}) = \begin{bmatrix} 1 - hr + \frac{1}{2}hr\ln(\frac{K}{\bar{u}^2}) & -\frac{hm}{2} \\ chm\bar{v} & 1 - dh + chm\bar{u} \end{bmatrix}.$$
 (2.1)

We apply the following results to examine the stability of the fixed points:

Lemma 2.1. [32] Let $P(q) = q^2 + Sq + T$ be the characteristic polynomial of $J(\bar{u}, \bar{v})$, and q_1, q_2 be the two roots of P(q) = 0.

- (*i*) If $|q_{1,2}| < 1$, then (\bar{u}, \bar{v}) is a sink (LAS),
- (ii) If $|q_{1,2}| > 1$, then (\bar{u}, \bar{v}) is a source,
- (iii) if $|q_1| < 1 \land |q_2| > 1$ (or $|q_1| > 1 \land |q_2| < 1$), then (\bar{u}, \bar{v}) is a saddle point (SP),
- (iv) if either $|q_1| = 1$ or $|q_2| = 1$, then (\bar{u}, \bar{v}) is a non-hyperbolic point (NHP).

Lemma 2.2. [32] Let $P(q) = q^2 + Sq + T$ and P(1) > 0. If q_1, q_2 are roots of P(q) = 0, then

- (i) $|q_{1,2}| < 1$ if $P(-1) > 0 \land T < 1$,
- (ii) $|q_1| < 1 \land |q_2| > 1$ (or $|q_1| > 1 \land |q_2| < 1$) if P(-1) < 0,
- (iii) $|q_1| > 1 \land |q_2| > 1$ if $P(-1) > 0 \land T > 1$,
- (iv) $q_1 = -1 \land |q_2| \neq 1$ if $P(-1) = 0 \land S \neq 0, 2$,
- (v) $q_1 \wedge q_2$ are complex, and $|q_{1,2}| = 1$ if $S^2 4T < 0 \wedge T = 1$.

The Jacobian matrix computed at the fixed point $P_1(\sqrt{K}, 0)$ is

$$J(P_1) = \begin{bmatrix} 1 - hr & -\frac{mh}{2} \\ 0 & 1 - dh + cmh\sqrt{K} \end{bmatrix}.$$

The following result describes the topological classification of $P_1(\sqrt{K}, 0)$:

Proposition 2.1. The fixed point $P_1(\sqrt{K}, 0)$ is

(i) LAS if

$$d > cm\sqrt{K}$$
 and $0 < h < \min\left\{\frac{2}{r}, \frac{2}{d-cm\sqrt{K}}\right\}$,

- (ii) a source if any of the following is true:
- (a) $d < cm\sqrt{K}$ and $h > \frac{2}{r}$, (b) $d > cm\sqrt{K}$ and $h > max\left\{\frac{2}{r}, \frac{2}{d-cm\sqrt{K}}\right\}$,
- (iii) a SP if any of the following is true:
- (a) $d < cm\sqrt{K}$ and $h < \frac{2}{r}$, (b) $d > cm\sqrt{K}$ and $\min\left\{\frac{2}{r}, \frac{2}{d-cm\sqrt{K}}\right\} < h < \max\left\{\frac{2}{r}, \frac{2}{d-cm\sqrt{K}}\right\}$,
- (iv) NHP if any of the following is true:
- (a) $h = \frac{2}{r}$, (b) $d = cm\sqrt{K}$,
- (c) $h = \frac{2}{d cm\sqrt{K}}, d \neq cm\sqrt{K}.$

It is evident that if $d = cm\sqrt{K}$, then one of the eigenvalues of $J(P_1)$ is 1. Consequently, a transcritical bifurcation may occur if the parameters change in a close neighborhood of Γ_1 .

$$\Gamma_1 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| d = d_1 = cm\sqrt{K}, h \neq \frac{2}{r} \right\}.$$

Moreover, if $h = \frac{2}{r}$ or $h = \frac{2}{d-cm\sqrt{K}}$, $d \neq cm\sqrt{K}$, then one of the eigenvalues of $J(P_1)$ is -1. So, a PD bifurcation may occur if the parameters change in a close neighborhood of Γ_2 or Γ_3 , where

$$\Gamma_2 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| h = h_1 = \frac{2}{r}, h \neq \frac{2}{d - cm\sqrt{K}}, d \neq cm\sqrt{K} \right\},\$$

and

$$\Gamma_3 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| h = h_{12} = \frac{2}{d - cm\sqrt{K}}, \ d \neq cm\sqrt{K}, \ h \neq \frac{2}{r} \right\}.$$

The Jacobian matrix at $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$ is

$$J(P_2) = \begin{bmatrix} 1 - hr + \frac{hr}{2} \ln(\frac{kc^2m^2}{d^2}) & -\frac{mh}{2} \\ \frac{dhr}{m} \ln(\frac{kc^2m^2}{d^2}) & 1 \end{bmatrix}.$$

The characteristic polynomial of $J(P_2)$ is

$$P(q) = q^2 + Sq + T,$$
 (2.2)

where

$$S = -2 + A_1 h$$
, $T = 1 - A_1 h + A_2 h^2$,

and

$$A_1 = r - \frac{r}{2} \ln\left(\frac{kc^2m^2}{d^2}\right), \ A_2 = \frac{dr}{2} \ln\left(\frac{kc^2m^2}{d^2}\right).$$

Through simple calculations, we obtain

$$P(0) = 1 - A_1h + A_2h^2,$$

$$P(1) = A_2h^2 > 0,$$

$$P(-1) = 4 - 2A_1h + A_2h^2.$$

Using the lemma (2.1), we acquire the local dynamics of the fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$.

Proposition 2.2. Assume that $k > \frac{d^2}{c^2m^2}$. The fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2}\ln(\frac{kc^2m^2}{d^2}))$ of the system (1.4) is

- (1) LAS if any of the following is true:
- (a) $A_1 > 0$, $A_1^2 4A_2 < 0$, and $0 < h < \frac{A_1}{A_2}$, (b) $A_1 > 0$, $A_1^2 - 4A_2 \ge 0$, and $0 < h < \frac{A_1 - \sqrt{A_1^2 - 4A_2}}{A_2}$,
- (2) source if any of the following is true:
- (a) $A_1 \leq 0$,
- (b) $A_1 > 0$, $A_1^2 4A_2 \le 0$, and $h > \frac{A_1}{A_2}$,
- (c) $A_1 > 0$, $A_1^2 4A_2 > 0$, and $h > \frac{A_1 + \sqrt{A_1^2 4A_2}}{A_2}$
- (3) SP if any of the following is true: $A_1 > 0$, $A_1^2 - 4A_2 > 0$, and $\frac{A_1 - \sqrt{A_1^2 - 4A_2}}{A_2} < h < \frac{A_1 + \sqrt{A_1^2 - 4A_2}}{A_2}$,
- (4) NHP if any of the following is true: (a) $A_1 > 0$, $A_1^2 - 4A_2 > 0$, and $h = \frac{A_1 \pm \sqrt{A_1^2 - 4A_2}}{A_2}$, (b) $A_1 > 0$, $A_1^2 - 4A_2 < 0$, and $h = \frac{A_1}{A_2}$.

The condition for $J(P_2)$ eigenvalues to be unit-modulus complex is established when $A_1 > 0$, $A_1^2 - 4A_2 < 0$, and $h = \frac{A_1}{A_2}$. Thus, the NS bifurcation occurs at point P_2 in the system (1.4) when the parameters are altered in the vicinity of Ω_1 .

$$\Omega_1 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| A_1 > 0, A_1^2 - 4A_2 < 0, h = h_3 = \frac{A_1}{A_2} \right\}.$$

Furthermore, if $A_1 > 0$, $A_1^2 - 4A_2 > 0$, and $h = \frac{A_1 \pm \sqrt{A_1^2 - 4A_2}}{A_2}$, one of the eigenvalues of the matrix $J(P_2)$ is -1, while the other eigenvalue λ meets the condition that $|\lambda| \neq 1$. If the parameters fluctuate slightly around Ω_2 or Ω_3 , a PD bifurcation may occur, where

$$\Omega_2 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| A_1 > 0, \ A_1^2 - 4A_2 > 0, \ h = h_{21} = \frac{A_1 + \sqrt{A_1^2 - 4A_2}}{A_2} \right\},\$$

and

$$\Omega_3 = \left\{ h, r, K, m, c, d \in \mathbb{R}^+ \middle| A_1 > 0, \ A_1^2 - 4A_2 > 0, \ h = h_2 = \frac{A_1 - \sqrt{A_1^2 - 4A_2}}{A_2} \right\}.$$

3. Local Bifurcation Analysis

The present section delves into various types of fixed-point bifurcations that could potentially manifest in the system (1.4). Bifurcations in predator-prey systems manifest as a result of alterations in the parameters. A slight modification in the parameter leads to a bifurcation. Bifurcations in predator-prey systems play a crucial role in forecasting the dynamics of wild populations and formulating viable approaches for sustainable management. Improper management of bifurcations can lead to disturbances in population dynamics and the consequent destruction of ecosystems. For a detailed bifurcation analysis, we recommend that readers refer to [31, 33–39].

3.1. Transcritical Bifurcation at $P_1(\sqrt{K}, 0)$.

This section discusses transcritical bifurcation at the boundary fixed point $P_1(\sqrt{K}, 0)$ using center manifold theory. Assuming that $(h, r, K, m, c, d) \in \Gamma_1$, and γ be small perturbation in d, the subsequent perturbation of the system (1.4) is taken into consideration:

$$\begin{cases} u_{n+1} = u_n + \frac{h}{2} (r u_n \ln(\frac{K}{u_n^2}) - m v_n), \\ v_{n+1} = v_n + h (c m u_n v_n - (d_1 + \gamma) v_n). \end{cases}$$
(3.1)

In order to translate fixed point $P_1(\sqrt{K}, 0)$ to (0, 0), we define the translation map as follows:

$$a_n = u_n - \sqrt{K}, \ b_n = v_n.$$

As a result of this translation map, the system (3.1) transforms to

$$\begin{bmatrix} a_{n+1} \\ b_{n+1} \end{bmatrix} = \begin{bmatrix} 1 - hr & -\frac{hm}{2} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} + \begin{bmatrix} -\frac{hr}{2\sqrt{K}}a_n^2 + \frac{hr}{6K}a_n^3 + O((|a_n| + |b_n| + |\gamma|)^4) \\ -h\gamma b_n + chma_nb_n + O((|a_n| + |b_n| + |\gamma|)^4) \end{bmatrix}.$$
 (3.2)

Let

$$T = \begin{bmatrix} -\frac{mr}{2} & 1\\ 1 & 0 \end{bmatrix}$$

Under the following transformation

$$\begin{bmatrix} a_n \\ b_n \end{bmatrix} = T \begin{bmatrix} e_n \\ f_n \end{bmatrix}, \qquad (3.3)$$

the system (3.2) becomes

$$\begin{bmatrix} e_{n+1} \\ f_{n+1} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1-hr \end{bmatrix} \begin{bmatrix} e_n \\ f_n \end{bmatrix} + \begin{bmatrix} -\frac{chm^2}{2r}e_n^2 + chme_nf_n - he_n\gamma + O((|e_n| + |f_n| + |\gamma|)^4) \\ F(e_n, f_n, \gamma) \end{bmatrix}, \quad (3.4)$$

where

$$F(e_n, f_n, \gamma) = -\frac{hm^2(2cm + \frac{r}{\sqrt{K}})}{8r^2}e_n^2 - \frac{hm^3}{48Kr^2}e_n^3 + \frac{hm}{2}(\frac{1}{\sqrt{K}} + \frac{cm}{r})e_nf_n + \frac{hm^2}{8Kr}e_n^2f_n - \frac{hr}{2\sqrt{K}}f_n^2 - \frac{hm}{4K}e_nf_n^2 + \frac{hr}{6K}f_n^3 - \frac{hm}{2r}e_n\gamma + O((|e_n| + |f_n| + |\gamma|)^4).$$

Subsequently, the center manifold is obtained for the system (3.4). This can be defined by the following:

$$W^{C} = \left\{ (e_{n}, f_{n}, \gamma) \in \mathbb{R}^{3} \middle| f_{n} = c_{1}e_{n}^{2} + c_{2}e_{n}\gamma + c_{3}\gamma^{2} + O((|e_{n}| + |\gamma|)^{3}) \right\},\$$

where

$$c_1 = -\frac{m^2(2cm + \frac{r}{\sqrt{\kappa}})}{8r^3}, \ c_2 = -\frac{m}{2r^2}, \ c_3 = 0.$$

Hence, the center manifold-restricted system (3.4) is

$$\tilde{F}: e_{n+1} = e_n - \frac{chm^2}{2r}e_n^2 - he_n\gamma - \frac{chm^3(2cm + \frac{r}{\sqrt{\kappa}})}{8r^3}e_n^3 - \frac{chm^2}{2r^2}e_n^2\gamma.$$

Since $\tilde{F}(0,0) = 0$, $\tilde{F}_{e_n}(0,0) = 1$, $\tilde{F}_{\gamma}(0,0) = 0$, $\tilde{F}_{e_ne_n}(0,0) = -\frac{chm^2}{r} \neq 0$, $\tilde{F}_{e_n\gamma}(0,0) = -h \neq 0$, the system (1.4) experiences transcritical bifurcation at the fixed point $P_1(\sqrt{K},0)$.

The following theorem provides the parametric conditions for the presence and direction of transcritical bifurcation for system (1.4) at its boundary fixed point $P_1(\sqrt{K}, 0)$.

Theorem 3.1. Assume that $(h, r, K, m, c, d) \in \Gamma_1$, then the system (1.4) experiences transcritical bifurcation at $P_1(\sqrt{K}, 0)$ when d varies in a small neighborhood of $d_1 = cm\sqrt{K}$. Furthermore, if $d > d_1$, then the system (1.4) has 2 fixed points, and if $d < d_1$, then the system (1.4) has only one fixed point $P_1(\sqrt{K}, 0)$.

3.2. **PD Bifurcation at** $P_1(\sqrt{K}, 0)$.

The PD bifurcation at the fixed point $P_1(\sqrt{K}, 0)$ in Γ_2 is discussed in this section. For the domain Γ_3 , a similar investigation can be done. Suppose that $(h, r, K, m, c, d) \in \Gamma_2$, and γ be minimal change in h, the subsequent perturbation of the system (1.4) is taken into consideration:

$$\begin{cases} u_{n+1} = u_n + \frac{h_1 + \gamma}{2} (r u_n \ln(\frac{K}{u_n^2}) - m v_n), \\ v_{n+1} = v_n + (h_1 + \gamma) (c m u_n v_n - d v_n). \end{cases}$$
(3.5)

In order to translate the fixed point $P_1(\sqrt{K}, 0)$ to (0, 0), we define the translation map as follows:

$$a_n = u_n - \sqrt{K}, \ b_n = v_n.$$

As a result of this translation map, the system (3.5) transforms to

$$\begin{bmatrix} a_{n+1} \\ b_{n+1} \end{bmatrix} = \begin{bmatrix} -1 & -\frac{m}{r} \\ 0 & \frac{-2d+2cm\sqrt{K}+r}{r} \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} + \begin{bmatrix} -\frac{m}{2}b_n\gamma - ra_n\gamma - \frac{1}{\sqrt{K}}a_n^2 - \frac{r}{2\sqrt{K}}a_n^2\gamma + \frac{1}{3K}a_n^3 + O((|a_n| + |b_n| + |\gamma|)^4) \\ (-d + cm\sqrt{K})b_n\gamma + \frac{2cm}{r}a_nb_n + cma_nb_n\gamma + O((|a_n| + |b_n| + |\gamma|)^4) \end{bmatrix}.$$
(3.6)

Let

$$T = \begin{bmatrix} 1 & -\frac{m}{2(-d+cm\sqrt{K}+r)} \\ 0 & 1 \end{bmatrix}.$$

Under the following transformation

$$\begin{bmatrix} a_n \\ b_n \end{bmatrix} = T \begin{bmatrix} e_n \\ f_n \end{bmatrix}, \tag{3.7}$$

system (3.6) becomes

$$\begin{bmatrix} e_{n+1} \\ f_{n+1} \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & \frac{-2d+2cm\sqrt{K}+r}{r} \end{bmatrix} \begin{bmatrix} e_n \\ f_n \end{bmatrix} + \begin{bmatrix} F(e_n, f_n, \gamma) \\ G(e_n, f_n, \gamma) \end{bmatrix},$$
(3.8)

where

$$\begin{split} F(e_n, f_n, \gamma) &= C_1 e_n^2 + C_2 e_n^3 + C_3 e_n f_n + C_4 e_n^2 f_n + C_5 f_n^2 + C_6 e_n f_n^2 + C_7 f_n^3 + C_8 e_n \gamma + C_9 e_n^2 \gamma + C_{10} e_n f_n \gamma \\ &+ O((|e_n| + |f_n| + |\gamma|)^4), \\ G(e_n, f_n, \gamma) &= \frac{2cm}{r} e_n f_n - \frac{cm^2}{r(-d + cm\sqrt{K} + r)} f_n^2 + (-d + cm\sqrt{K}) f_n \gamma + cme_n f_n \gamma + O((|e_n| + |f_n| + |\gamma|)^4), \\ C_1 &= -\frac{1}{\sqrt{K}}, \ C_2 &= \frac{1}{3K}, \ C_3 &= \frac{m(cm\sqrt{K} + r)}{r\sqrt{K}(-d + cm\sqrt{K} + r)}, \ C_4 &= \frac{m}{2dK - 2cmK^{3/2} - 2Kr}, \\ C_5 &= -\frac{m^2(2cm\sqrt{K} + r)}{4r\sqrt{K}(-d + cm\sqrt{K} + r)^2}, \ C_6 &= \frac{m^2}{4K(-d + cm\sqrt{K} + r)^2}, \ C_7 &= -\frac{m^3}{24K(-d + cm\sqrt{K} + r)^3}, \\ C_8 &= -r, \ C_9 &= \frac{-r}{\sqrt{K}}, \ C_{10} &= \frac{m(cm\sqrt{K} + r)}{2\sqrt{K}(-d + cm\sqrt{K} + r)}. \end{split}$$

Subsequently, the center manifold W^C for the system (3.8) is computed and characterized as follows:

$$W^{C} = \left\{ (e_{n}, f_{n}, \gamma) \in \mathbb{R}^{3} \middle| f_{n} = c_{1}e_{n}^{2} + c_{2}e_{n}\gamma + c_{3}\gamma^{2} + O((|e_{n}| + |\gamma|)^{3}) \right\},\$$

where

$$c_1 = 0, \ c_2 = 0, \ c_3 = 0.$$

System (3.8) when confined to the center manifold is expressed as follows:

$$\tilde{F}: e_{n+1} = -e_n - \frac{1}{\sqrt{K}}e_n^2 + \frac{1}{3K}e_n^3 - re_n\gamma - \frac{r}{\sqrt{K}}e_n^2\gamma.$$
(3.9)

For map (3.9) to experience PD bifurcation, we need the value of the following two expressions to be non-zero:

$$l_{1} = \tilde{F}_{\gamma}\tilde{F}_{e_{n}e_{n}} + 2\tilde{F}_{e_{n}\gamma}\Big|_{(0,0)}, \ l_{2} = \frac{1}{2}(\tilde{F}_{e_{n}e_{n}})^{2} + \frac{1}{3}\tilde{F}_{e_{n}e_{n}e_{n}}$$

From simple computations, we obtain

$$l_1 = -2r, \ l_2 = \frac{8}{3K}.$$

The following theorem provides the parametric conditions for the presence and direction of PD bifurcation for system (1.4) at its boundary fixed point $P_1(\sqrt{K}, 0)$.

Theorem 3.2. Assume that $(h, r, K, m, c, d) \in \Gamma_2$, then the system (1.4) experiences PD bifurcation at $P_1(\sqrt{K}, 0)$ when h varies in a small neighborhood of $h_1 = \frac{2}{r}$. Furthermore, the period-2 points that bifurcate from $P_1(\sqrt{K}, 0)$ are stable.

3.3. **PD Bifurcation at** $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$. The PD bifurcation at the fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$ for the domain Ω_3 is discussed in this section. For the domain Ω_2 , a similar investigation can be done. Suppose that $(h, r, K, m, c, d) \in$ Ω_3 , and γ be minimal change in h, the subsequent perturbation of the system (1.4) is taken into consideration:

$$\begin{cases} u_{n+1} = u_n + \frac{h_2 + \gamma}{2} (r u_n \ln(\frac{\kappa}{u_n^2}) - m v_n), \\ v_{n+1} = v_n + (h_2 + \gamma) (c m u_n v_n - d v_n). \end{cases}$$
(3.10)

In order to translate fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2}\ln(\frac{kc^2m^2}{d^2}))$ to (0,0), we define the translation map as follows:

$$a_n = u_n - \frac{d}{cm}, \ b_n = v_n - \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}).$$

As a result of this translation map, the system (3.10) transforms to

$$\begin{bmatrix} a_{n+1} \\ b_{n+1} \end{bmatrix} = \begin{bmatrix} 1 - A_1 h_2 & -\frac{h_2 m}{2} \\ \frac{2h_2 A_2}{m} & 1 \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} + \begin{bmatrix} F(a_n, b_n, \gamma) \\ G(a_n, b_n, \gamma) \end{bmatrix},$$
(3.11)

where

$$F(a_{n}, b_{n}, \gamma) = -\frac{m}{2}b_{n}\gamma - \frac{ch_{2}mr}{2d}a_{n}^{2} - \frac{cmr}{2d}a_{n}^{2}\gamma + \frac{c^{2}h_{2}m^{2}r}{6d^{2}}a_{n}^{3} - A_{1}a_{n}\gamma + O((|a_{n}| + |b_{n}| + |\gamma|)^{4}),$$

$$G(a_{n}, b_{n}, \gamma) = ch_{2}ma_{n}b_{n} + cma_{n}b_{n}\gamma + \frac{2A_{2}}{m}a_{n}\gamma + O((|a_{n}| + |b_{n}| + |\gamma|)^{4}).$$
For $h_{2} = \frac{A_{1} - \sqrt{A_{1}^{2} - 4A_{2}}}{A_{2}}$, the eigenvalues of $J(P_{2})$ are $\lambda_{1} = -1$ and $\lambda_{2} = 3 - A_{1}h_{2}.$
Let

$$T = \begin{bmatrix} \frac{m}{-A_1 + \sqrt{A_1^2 - 4A_2}} & -\frac{-A_1^2 m + A_1 m \sqrt{A_1^2 - 4A_2} + 2A_2 m}{2A_2(-A_1 + \sqrt{A_1^2 - 4A_2})} \\ 1 & 1 \end{bmatrix}.$$

Under the following transformation

$$\begin{bmatrix} a_n \\ b_n \end{bmatrix} = T \begin{bmatrix} e_n \\ f_n \end{bmatrix}.$$
 (3.12)

System (3.11) becomes

$$\begin{bmatrix} e_{n+1} \\ f_{n+1} \end{bmatrix} = \begin{bmatrix} -1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \begin{bmatrix} e_n \\ f_n \end{bmatrix} + \begin{bmatrix} F(e_n, f_n, \gamma) \\ G(e_n, f_n, \gamma) \end{bmatrix},$$
(3.13)

where

$$\lambda_{2} = \frac{2A_{1}^{4} - 2A_{1}^{3}\sqrt{A_{1}^{2} - 4A_{2}} - 11A_{1}^{2}A_{2} + 7A_{1}A_{2}\sqrt{A_{1}^{2} - 4A_{2}} + 12A_{2}^{2}}{A_{2}(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},$$

$$F(e_{n}, f_{n}, \gamma) = D_{1}e_{n}^{2} + D_{2}e_{n}^{3} + D_{3}e_{n}f_{n} + D_{4}e_{n}^{2}f_{n} + D_{5}f_{n}^{2} + D_{6}e_{n}f_{n}^{2} + D_{7}f_{n}^{3} + D_{8}e_{n}\gamma + D_{9}e_{n}^{2}\gamma + D_{10}e_{n}f_{n}\gamma + O((|e_{n}| + |f_{n}| + |\gamma|)^{4}),$$

 $G(e_n, f_n, \gamma) = E_1 e_n^2 + E_2 e_n^3 + E_3 e_n f_n + E_4 e_n^2 f_n + E_5 f_n^2 + E_6 e_n f_n^2 + E_7 f_n^3 + E_8 e_n^2 \gamma + E_9 f_n \gamma + E_{10} e_n f_n \gamma + O((|e_n| + |f_n| + |\gamma|)^4),$

$$\begin{split} D_1 &= -\frac{cm^2(-A_1^2d + A_1d\sqrt{A_1^2 - 4A_2} + A_2(2d - r))}{A_2d(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_2 &= -\frac{c^2m^4r}{3d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_3 &= -\frac{cm^2(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 2A_2)(A_1^2d - A_1d\sqrt{A_1^2 - 4A_2} + 2A_2r)}{2A_2^2d(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_4 &= \frac{c^2m^4r(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 2A_2)}{2A_2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_5 &= \frac{cm^2(2d + r)(A_1^2 - A_1\sqrt{A_1^2 - 4A_2} - 2A_2)^2}{4A_2^2d(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_6 &= -\frac{c^2m^4r(A_1^2 - A_1\sqrt{A_1^2 - 4A_2} - 2A_2)^2}{4A_2^2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_6 &= -\frac{c^2m^4r(A_1^2 - A_1\sqrt{A_1^2 - 4A_2} - 2A_2)^2}{4A_2^2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_6 &= -\frac{c^2m^4r(A_1^2 - A_1\sqrt{A_1^2 - 4A_2} - 2A_2)^2}{4A_2^2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_6 &= -\frac{c^2m^4r(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} - 2A_2)^2}{4A_2^2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_7 &= \frac{c^2m^4r(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}{2A_2d^2(-A_1 + \sqrt{A_1^2 - 4A_2})(-A_1^2 + A_1\sqrt{A_1^2 - 4A_2} + 4A_2)}, \\ D_8 &= \frac{2A_2}{-A_1 + \sqrt{A_1^2 - 4A_2}}, \end{split}$$

$$\begin{split} D_{0} &= \frac{2cm^{2}(-A_{1}^{2}d + A_{1}d\sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + A_{2}(2d - r))}{d(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 2A_{2})(-A_{1}^{2}d + A_{1}d\sqrt{A_{1}^{2} - 4A_{2}} - 2A_{2}r)}{2A_{2}d(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 2A_{2})},\\ E_{1} &= -\frac{cm^{2}(2d + r)}{d(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{2} &= \frac{c^{2}m^{4}r}{3d^{2}(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{3} &= \frac{cm^{2}(2A_{2}r - A_{1}^{2}(d + r) + A_{1}(d + r)\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}{A_{2}d(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{4} &= \frac{c^{2}m^{4}r(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 2A_{2})}{2A_{2}d^{2}(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 2A_{2})},\\ E_{5} &= \frac{cm^{2}(-2A_{1}^{2}A_{2}(d - 2r) + 2A_{1}A_{2}(d - r)\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}{2A_{2}d(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 2A_{2})},\\ E_{6} &= \frac{c^{2}m^{4}r(A_{1}^{2} - A_{1}\sqrt{A_{1}^{2} - 4A_{2}} - 2A_{2})^{2}}{AA_{2}^{2}d^{2}(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{7} &= -\frac{c^{2}m^{4}r(A_{1}^{2} - A_{1}\sqrt{A_{1}^{2} - 4A_{2}} - 2A_{2})^{2}}{2A_{4}A_{2}^{2}d^{2}(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{8} &= \frac{2(-A_{1}^{4} + A_{1}^{3}\sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}}{d(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{9} &= \frac{2(-A_{1}^{4} + A_{1}^{3}\sqrt{A_{1}^{2} - 4A_{2}} + 5A_{1}^{3}A_{2} - 3A_{1}A_{2}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}{(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})},\\ E_{10} &= -\frac{cm^{2}(2A_{2}r - A_{1}^{2}(d + r) + A_{1}(d + r)\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}{d(-A_{1} + \sqrt{A_{1}^{2} - 4A_{2}})(-A_{1}^{2} + A_{1}\sqrt{A_{1}^{2} - 4A_{2}} + 4A_{2})}.\\ \end{array}$$

Next, we derive the center manifold for (3.13), which may be characterized as follows:

$$W^{C} = \left\{ (e_{n}, f_{n}, \gamma) \in \mathbb{R}^{3} \middle| f_{n} = c_{1}e_{n}^{2} + c_{2}e_{n}\gamma + c_{3}\gamma^{2} + O((|e_{n}| + |\gamma|)^{3}) \right\},\$$

where

$$c_1 = \frac{A_2 c m^2 (2d+r)}{2d(A_1^4 + A_1^3 \sqrt{A_1^2 - 4A_2} - 5A_1^2 A_2 - 3A_1 A_2 \sqrt{A_1^2 - 4A_2} + 4A_2^2)}, \ c_2 = 0, \ c_3 = 0.$$

System (3.13) restricted to the center manifold is

$$\tilde{F}: e_{n+1} = -e_n + D_1 e_n^2 + D_2 e_n^3 + D_8 e_n \gamma + D_9 e_n^2 \gamma + \frac{D_3 E_1}{1 - \lambda_2} e_n^3.$$
(3.14)

For map (3.14) to experience PD bifurcation, we need the value of the following two expressions to be non-zero:

$$l_{1} = \tilde{F}_{\gamma}\tilde{F}_{e_{n}e_{n}} + 2\tilde{F}_{e_{n}\gamma}\Big|_{(0,0)}, \ l_{2} = \frac{1}{2}(\tilde{F}_{e_{n}e_{n}})^{2} + \frac{1}{3}\tilde{F}_{e_{n}e_{n}e_{n}}.$$

From simple computations, we obtain

$$l_1 = 2D_8, \ l_2 = 2(D_1^2 + D_2 + \frac{D_3E_1}{1 - \lambda_2}).$$

The following result is drawn from the calculations mentioned above:

Theorem 3.3. Suppose that $(h, r, K, m, c, d) \in \Omega_3$. The system (1.4) undergoes PD bifurcation at the fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$ if $l_2 \neq 0$ and h differs in a small neighborhood of $h_2 = \frac{A_1 - \sqrt{A_1^2 - 4A_2}}{A_2}$. Moreover, if $l_2 > 0$ (respectively $l_2 < 0$), then the period-2 orbits that bifurcate from $P_2(\frac{d}{cm}, \frac{rd}{d^2} \ln(\frac{kc^2m^2}{d^2}))$ are stable (respectively, unstable).

3.4. **NS Bifurcation at** $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$.

This section discusses the NS bifurcation at the fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$ for the domain Ω_1 . Suppose that $(h, r, K, m, c, d) \in \Omega_1$, and γ be minimal change in h_3 , the subsequent perturbation of the system (1.4) is taken into consideration:

$$\begin{cases} u_{n+1} = u_n + \frac{h_3 + \gamma}{2} (r u_n \ln(\frac{K}{u_n^2}) - m v_n), \\ v_{n+1} = v_n + (h_3 + \gamma) (c m u_n v_n - d v_n). \end{cases}$$
(3.15)

We define $a_n = u_n - \frac{d}{cm}$, $b_n = v_n - \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2})$, to translate fixed point $P_2(\frac{d}{cm}, \frac{rd}{cm^2} \ln(\frac{kc^2m^2}{d^2}))$ to (0,0). As a result of this translation map, the system (3.15) transforms to

$$\begin{bmatrix} a_{n+1} \\ b_{n+1} \end{bmatrix} = \begin{bmatrix} 1 - \frac{A_1(A_1 + A_2\gamma)}{A_2} & -\frac{m}{2}(\frac{A_1}{A_2} + \gamma) \\ \frac{2A_1 + 2A_2\gamma}{m} & 1 \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} + \begin{bmatrix} F(a_n, b_n) \\ G(a_n, b_n) \end{bmatrix},$$
(3.16)

where

$$F(a_n, b_n) = -\frac{cmr(\frac{A_1}{A_2} + \gamma)}{2d}a_n^2 + \frac{c^2m^2r(\frac{A_1}{A_2} + \gamma)}{6d^2}a_n^3 + O((|a_n| + |b_n| + |\gamma|)^4),$$

$$G(a_n, b_n) = cm(\frac{A_1}{A_2} + \gamma)a_nb_n + O((|a_n| + |b_n| + |\gamma|)^4).$$

The characteristic equation corresponding to the linearized part of the system (3.16) at the origin is

$$\lambda^2 - \alpha(\gamma)\lambda + \beta(\gamma) = 0, \qquad (3.17)$$

where

$$\alpha(\gamma) = 2 - \frac{A_1^2}{A_2} - A_1\gamma,$$
$$\beta(\gamma) = 1 + A_1\gamma + A_2\gamma^2.$$

The equation (3.17) roots are complex that have the property $|\lambda_{1,2}| = 1$, which are given by

$$\lambda_{1,2} = \frac{\alpha(\gamma) \pm i\sqrt{4\beta(\gamma) - \alpha^2(\gamma)}}{2}$$

By computations, we obtain

$$|\lambda_1| = |\lambda_2| = \sqrt{\beta(\gamma)},$$

and

$$\left(\frac{d|\lambda_1|}{d\gamma}\right)_{\gamma=0} = \left(\frac{d|\lambda_2|}{d\gamma}\right)_{\gamma=0} = \frac{A_1}{2} > 0.$$

Moreover, it is required that λ_1^i , $\lambda_2^i \neq 1$ for i = 1, 2, 3, 4 at $\gamma = 0$, which is equivalent to $\alpha(0) \neq \pm 2, 0, 1$. Since $A_1 > 0$, $A_1^2 - 4A_2 < 0$ and $\alpha(0) = 2 - \frac{A_1^2}{A_2}$, therefore $\alpha(0) \neq \pm 2$. We only need to require that $\alpha(0) \neq 0, 1$, which leads to $A_1^2 \neq 2A_2, A_2$.

The canonical form at $\gamma = 0$ of the linear part in (3.16) is achieved through the utilization of the subsequent transformation:

$$\begin{bmatrix} a_n \\ b_n \end{bmatrix} = \begin{bmatrix} -\frac{A_1m}{2A_2} & 0 \\ \frac{A_1^2}{2A_2} & -\frac{A_1\sqrt{4A_2-A_1^2}}{2A_2} \end{bmatrix} \begin{bmatrix} e_n \\ f_n \end{bmatrix}.$$
 (3.18)

Upon application of the mapping (3.18), the system (3.16) transforms as follows:

$$\begin{bmatrix} e_{n+1} \\ f_{n+1} \end{bmatrix} = \begin{bmatrix} \mu & -\nu \\ \nu & \mu \end{bmatrix} \begin{bmatrix} e_n \\ f_n \end{bmatrix} + \begin{bmatrix} F(e_n, f_n) \\ G(e_n, f_n) \end{bmatrix},$$
(3.19)

where

$$\mu = 1 - \frac{A_1^2}{2A_2}, \ \nu = \frac{A_1\sqrt{4A_2 - A_1^2}}{2A_2},$$

$$F(e_n, f_n) = \frac{A_1^2 cm^2 r}{4A_2^2 d} e_n^2 + \frac{A_1^3 c^2 m^4 r}{24A_2^3 d^2} e_n^3 + O((|e_n| + |f_n| + |\gamma|)^4),$$

$$G(e_n, f_n) = \frac{A_1^3 cm^2 (2d + r)}{4A_2^2 d\sqrt{4A_2 - A_1^2}} e_n^2 + \frac{A_1^4 c^2 m^4 r}{24A_2^3 d^2 \sqrt{4A_2 - A_1^2}} e_n^3 - \frac{A_1^2 cm^2}{2A_2^2} e_n f_n + O((|e_n| + |f_n| + |\gamma|)^4).$$

The below-mentioned number L explains how the invariant curve appears in a system going through the NS bifurcation.

$$L = \left(\left[-Re\left(\frac{(1-2\lambda_1)\lambda_2^2}{1-\lambda_1} \eta_{20} \eta_{11} \right) - \frac{1}{2} |\eta_{11}|^2 - |\eta_{02}|^2 + Re(\lambda_2 \eta_{21}) \right] \right)_{\gamma=0}$$

where

$$\begin{split} \eta_{20} &= \frac{1}{8} \left[F_{e_n e_n} - F_{f_n f_n} + 2G_{e_n f_n} + i(G_{e_n e_n} - G_{f_n f_n} - 2F_{e_n f_n}) \right], \\ \eta_{11} &= \frac{1}{4} \left[F_{e_n e_n} + F_{f_n f_n} + i(G_{e_n e_n} + G_{f_n f_n}) \right], \\ \eta_{02} &= \frac{1}{8} \left[F_{e_n e_n} - F_{f_n f_n} - 2G_{e_n f_n} + i(G_{e_n e_n} - G_{f_n f_n} + 2F_{e_n f_n}) \right], \\ \eta_{21} &= \frac{1}{16} \left[F_{e_n e_n e_n} + F_{e_n f_n f_n} + G_{e_n e_n f_n} + G_{f_n f_n f_n} + i(G_{e_n e_n e_n} + G_{e_n f_n f_n} - F_{e_n e_n f_n} - F_{f_n f_n f_n}) \right]. \end{split}$$

From the calculations shown above, we can obtain the following result:

Theorem 3.4. Consider $(h, r, K, m, c, d) \in \Omega_1$ and $A_1^2 \neq 2A_2, A_2$. When the parameter h differs in a small neighborhood of $h_3 = \frac{A_1}{A_2}$, the system (1.4) undergoes NS bifurcation at the fixed point P_2 if $L \neq 0$. Additionally, if L > 0, a repelling invariant closed curve bifurcates from P_2 , whereas if L < 0, an attracting invariant closed curve bifurcates from P_2 .

4. Numerical examples

This section will strengthen our theoretical investigations of the system by presenting some numerical examples that illustrate its numerous qualitative properties.

4.1. PD bifurcation of the system (1.4) at P_2 by using h as bifurcation parameter.

Consider the following parametric values:

$$c = 0.9, m = 0.5, d = 0.15, r = 1.5, k = 0.12,$$

and initial conditions $u_0 = 0.5$, $v_0 = 0.05$. For these values, the bifurcation value is h = 1.39251, and the positive fixed point is obtained as $P_2(0.333333, 0.076961)$. The eigenvalues of $J(P_2)$ are obtained as $\lambda_1 = -1$, $\lambda_2 = 0.991606$, confirming that the system (1.4) undergoes period doubling bifurcation at $P_2(0.333333, 0.076961)$ as h passes through $h_0 = 1.39251$. Figures (1a,1b) depict bifurcation figures for $h \in [1.1, 1.7]$. The MLE is plotted in figure (1c).

The fixed point is a LAS for these parametric values if 0 < h < 1.39251. Figures (1d,1e,1f) depict phase portraits of the system (1.4) for distinct values of *h*. From the figures, the $P_2(0.333333, 0.076961)$ is a LAS for 0 < h < 1.39251, but fails to retain stability at h = 1.39251, where the system (1.4) experiences PD bifurcation.

4.2. NS bifurcation of the system (1.4) at P_2 by using h as bifurcation parameter.

Consider the following parametric values:

$$c = 0.9, m = 0.5, d = 0.15, r = 1.5, k = 0.5,$$

and initial conditions $u_0 = 0.4$, $v_0 = 1.4$. For these values, the bifurcation value is h = 2.19813, and the positive fixed point is obtained as $P_2(0.333333, 1.50408)$. The eigenvalues of $J(P_2)$ are computed as $\lambda_1 = 0.591212 - 0.806516i$, $\lambda_2 = 0.591212 + 0.806516i$ with $|\lambda_{1,2}| = 1$, confirming that the system (1.4) is going through NS bifurcation at $P_2(0.333333, 1.50408)$ as h passes through $h_0 = 2.19813$. Figures (2a,2b) depict bifurcation diagrams for $h \in [2.1, 2.8]$. The MLE is plotted in figure 2c.

The fixed point P_2 is a LAS for these parametric values if h < 2.19813. Figures (2d-2k) depict phase portraits of the system (1.4) for some values of h. From the figures, P_2 is a LAS for h < 2.19813 but loses stability at h = 2.19813, where the system (1.4) goes through NS bifurcation. A smooth curve that is invariant emerges for values of $h \ge 2.19813$, with its radius increasing proportionally to the increase in h. The sudden disappearance of the invariant curve and subsequent emergence of a



Figure 1. Bifurcation diagrams, MLE graph, phase portraits for c = 0.9, m = 0.5, d = 0.15, r = 1.5, k = 0.12, $u_0 = 0.5$, $v_0 = 0.05$.

periodic orbit are observed upon increasing the value of h. Subsequently, the periodic orbit is replaced by an invariant curve. The phenomenon persists until a strange attractor appears. See Figures (2l-2o).

4.3. NS bifurcation of the system (1.4) at P_2 by using d as bifurcation parameter.

Consider the following parametric values:

$$c = 0.9, m = 0.5, r = 1.5, k = 0.5, h = 2.19813$$

and initial conditions $u_0 = 0.4$, $v_0 = 1.4$. For these values, the bifurcation value is d = 0.15, and the positive fixed point is obtained as $P_2(0.333333, 1.50408)$. The eigenvalues of $J(P_2)$ are $\lambda_1 = 0.591212 - 0.806518i$, $\lambda_2 = 0.591212 + 0.806518i$ with $|\lambda_{1,2}| = 1$, indicating that the system (1.4) is experiencing NS bifurcation at $P_2(0.333333, 1.50408)$ as the bifurcation parameter d passes through $d_0 = 0.15$. Figures (3a,3b) depict bifurcation diagrams for $d \in [0.14, 0.17]$. The MLE is plotted in figure 3c. The bifurcation diagrams show that the positive fixed point is stable if the death rate of the predator is greater than a threshold value. Biologically, it means that to prevent the predator population from growing uncontrollably and stability of the positive fixed point, the predator's death rate should be greater than the threshold value.

4.4. Sensitive dependence on initial conditions.

Sensitive dependency on initial conditions implies that small changes in the initial populations of predator and prey may result in drastically different long-term population dynamics. This phenomenon highlights the importance of accurate and precise measurements of the initial conditions in ecological investigations and the requirement for strong modeling tools that can account for data uncertainty and variability. It also highlights the difficulties in forecasting the long-term dynamics of ecological systems since even minor errors in the initial conditions may lead to considerable errors in predictions.

Figure (4) depicts 2 perturbed trajectories in blue and red colors to highlight the sensitivity of the system (1.4) to initial conditions. The two trajectories are initially overlapping and identical, but the divergence between them grows fast after a few repetitions. Figure (4) plots the u- and v-coordinates of the two trajectories for the system (1.4) against the number of iterations, revealing a reactive reliance on the initial conditions. The initial perturbation of 2 trajectories is 0.0001. The 2 trajectories with initial points (u_0 , v_0) = (0.4, 1.4) (trajectory in blue color) and (u_0 , v_0) = (0.4001, 1.4001) (trajectory in red color) are calculated and plotted in figure (4), respectively. The trajectories of the system (1.4) are depicted in figure (4) to be sensitively dependent on the initial conditions, i.e., complex dynamic behavior happens with an initial perturbation.



Figure 2. Bifurcation diagrams, MLE graph, phase portraits for c = 0.9, m = 0.5, d = 0.15, r = 1.5, k = 0.5, $u_0 = 0.4$, $v_0 = 1.4$.



Figure 3. Bifurcation diagrams, MLE graph for c = 0.9, m = 0.5, r = 1.5, k = 0.5, h = 2.19813, $u_0 = 0.4$, $v_0 = 1.4$.



Figure 4. Sensitive dependence on initial conditions of the system (1.4)

5. Conclusion

The nonlinear dynamics of a novel discrete-time predator-prey system with square root functional response and Gompertz growth of prey, generated using the forward Euler discretization method, were examined in this study. The system has a boundary fixed point $(\sqrt{K}, 0)$ which always exists, and an interior fixed point $(\frac{d}{cm}, \frac{rd}{cm^2}ln(\frac{kc^2m^2}{d^2}))$ which exists only if the carrying capacity of prey is greater than a threshold value. Using bifurcation theory and the center manifold theorem, it is shown that the system's fixed points have transcritical bifurcation, PD bifurcation, and NS bifurcation. The less integral step size *h* can stabilize the positive fixed point. However, the large integral step size can destabilize the positive fixed point, resulting in more complex dynamical behaviors, as depicted in the figures. Moreover, we see that if the death rate of the predator is greater than a threshold value, then the positive fixed point is stable; otherwise, it is unstable. This threshold value shows the lowest death rate necessary to prevent the predator population from exploding and to maintain a stable equilibrium between predator and prey populations. Moreover, it is demonstrated that the trajectories of the system (1.4) are sensitive to the initial values. Slight variations in initial conditions can significantly affect the long-term dynamics of ecological systems. The sensitive dependence on initial conditions

emphasizes the importance of accurate and precise measurements and the need for robust modeling techniques in ecological studies.

Conflicts of Interest: The authors declare that there are no conflicts of interest regarding the publication of this paper.

References

- A.J. Lotka, Elements of Physical Biology, Science Progress in the Twentieth Century (1919-1933), Vol. 21, pp. 341–343, Williams & Wilkins Company, Philadelphia, (1926).
- [2] V. Volterra, Fluctuations in the abundance of a species considered mathematically¹, Nature. 118 (1926), 558-560. https://doi.org/10.1038/118558a0.
- [3] C.S. Holling, Some Characteristics of Simple Types of Predation and Parasitism, Canadian Entomol. 91 (1959), 385–398.
- P.H. Crowley, E.K. Martin, Functional Responses and Interference within and between Year Classes of a Dragonfly Population, J. North Amer. Benthol. Soc. 8 (1989), 211–221. https://doi.org/10.2307/1467324.
- [5] J.R. Beddington, Mutual Interference Between Parasites or Predators and its Effect on Searching Efficiency, J. Animal Ecol. 44 (1975), 331–340. https://doi.org/10.2307/3866.
- [6] D.L. DeAngelis, R.A. Goldstein, R.V. O'Neill, A Model for Tropic Interaction, Ecology. 56 (1975), 881–892. https://doi.org/10.2307/1936298.
- M.F. Elettreby, A. Khawagi, T. Nabil, Dynamics of a Discrete Prey-Predator Model with Mixed Functional Response, Int. J. Bifurcation Chaos. 29 (2019), 1950199. https://doi.org/10.1142/s0218127419501992.
- [8] S.M.S. Rana, U. Kulsum, Bifurcation Analysis and Chaos Control in a Discrete-Time Predator-Prey System of Leslie Type with Simplified Holling Type IV Functional Response, Discr. Dyn. Nat. Soc. 2017 (2017), 9705985. https://doi.org/10.1155/2017/9705985.
- X. Chen, X. Zhang, Dynamics of the Predator-Prey Model With the Sigmoid Functional Response, Stud. Appl. Math. 147 (2021), 300-318. https://doi.org/10.1111/sapm.12382.
- [10] C. Arancibia-Ibarra, P. Aguirre, J. Flores, P. van Heijster, Bifurcation Analysis of a Predator-Prey Model With Predator Intraspecific Interactions and Ratio-Dependent Functional Response, Appl. Math. Comput. 402 (2021), 126152. https://doi.org/10.1016/j.amc.2021.126152.
- [11] S.M.S. Rana, Bifurcation and Complex Dynamics of a Discrete-Time Predator-Prey System With Simplified Monod-Haldane Functional Response, Adv. Differ. Equ. 2015 (2015), 345. https://doi.org/10.1186/ s13662-015-0680-7.
- H. Baek, Complex Dynamics of a Discrete-Time Predator-Prey System with Ivlev Functional Response, Math. Probl. Eng. 2018 (2018), 8635937. https://doi.org/10.1155/2018/8635937.
- [13] V. Ajraldi, M. Pittavino, E. Venturino, Modeling Herd Behavior in Population Systems, Nonlinear Anal.: Real World Appl. 12 (2011), 2319–2338. https://doi.org/10.1016/j.nonrwa.2011.02.002.
- [14] P.A. Braza, Predator-Prey Dynamics With Square Root Functional Responses, Nonlinear Anal.: Real World Appl. 13 (2012), 1837–1843. https://doi.org/10.1016/j.nonrwa.2011.12.014.
- [15] S.M. Salman, A.M. Yousef, A.A. Elsadany, Stability, Bifurcation Analysis and Chaos Control of a Discrete Predator-Prey System With Square Root Functional Response, Chaos Solitons Fractals. 93 (2016), 20–31. https://doi. org/10.1016/j.chaos.2016.09.020.
- [16] M. Berkal, M.B. Almatrafi, Bifurcation and Stability of Two-Dimensional Activator–Inhibitor Model with Fractional-Order Derivative, Fractal Fract. 7 (2023), 344. https://doi.org/10.3390/fractalfract7050344.

- [17] M.B. Almatrafi, Solitary Wave Solutions to a Fractional Model Using the Improved Modified Extended Tanh-Function Method, Fractal Fract. 7 (2023), 252. https://doi.org/10.3390/fractalfract7030252.
- [18] A.Q. Khan, M.B. Almatrafi, Two-Dimensional Discrete-Time Laser Model With Chaos and Bifurcations, AIMS Math. 8 (2023), 6804–6828. https://doi.org/10.3934/math.2023346.
- [19] A.Q. Khan, S.A.H. Bukhari, M.B. Almatrafi, Global Dynamics, Neimark-Sacker Bifurcation and Hybrid Control in a Leslie's Prey-Predator Model, Alexandria Eng. J. 61 (2022), 11391–11404. https://doi.org/10.1016/j.aej. 2022.04.042.
- [20] M.B. Almatrafi, Abundant Traveling Wave and Numerical Solutions for Novikov-Veselov System With Their Stability and Accuracy, Appl. Anal. 102 (2022), 2389–2402. https://doi.org/10.1080/00036811.2022.2027381.
- [21] M.G. Mortuja, M.K. Chaube, S. Kumar, Dynamic Analysis of a Predator-Prey System With Nonlinear Prey Harvesting and Square Root Functional Response, Chaos Solitons Fractals. 148 (2021), 111071. https://doi.org/ 10.1016/j.chaos.2021.111071.
- [22] P. Chakraborty, U. Ghosh, S. Sarkar, Stability and Bifurcation Analysis of a Discrete Prey-predator Model With Square-Root Functional Response and Optimal Harvesting, J. Biol. Syst. 28 (2020), 91-110. https://doi.org/ 10.1142/s0218339020500047.
- [23] P. Panja, Combine Effects of Square Root Functional Response and Prey Refuge on Predator-Prey Dynamics, Int. J. Model. Simul. 41 (2020), 426-433. https://doi.org/10.1080/02286203.2020.1772615.
- [24] B. Gompertz, On the Nature of the Function Expressive of the Law of Human Mortality, and on a New Mode of Determining the Value of Life Contingencies, Phil. Trans. R. Soc. Lond. 115 (1825), 513–583. https://doi.org/ 10.1098/rstl.1825.0026.
- [25] P.A. Naik, Z. Eskandari, M. Yavuz, J. Zu, Complex Dynamics of a Discrete-Time Bazykin-Berezovskaya Prey-Predator Model With a Strong Allee Effect, J. Comput. Appl. Math. 413 (2022), 114401. https://doi.org/10. 1016/j.cam.2022.114401.
- [26] S.M.S. Rana, Dynamics and Chaos Control in a Discrete-Time Ratio-Dependent Holling-Tanner Model, J. Egypt Math. Soc. 27 (2019), 48. https://doi.org/10.1186/s42787-019-0055-4.
- [27] M. Zhao, C. Li, J. Wang, Complex Dynamic Behaviors of a Discrete-Time Predator-Prey System, J. Appl. Anal. Comput. 7 (2017), 478–500. https://doi.org/10.11948/2017030.
- [28] K.S.N. Al-Basyouni, A.Q. Khan, Bifurcation Analysis of a Discrete-Time Chemostat Model, Math. Probl. Eng. 2023 (2023), 7518261. https://doi.org/10.1155/2023/7518261.
- [29] B. Hong, C. Zhang, Neimark-Sacker Bifurcation of a Discrete-Time Predator-Prey Model with Prey Refuge Effect, Mathematics. 11 (2023), 1399. https://doi.org/10.3390/math11061399.
- [30] Q. Din, Complex Dynamical Behavior and Control of a Discrete Ecological Model, J. Vib. Control. (2022). https: //doi.org/10.1177/10775463221133427.
- [31] A. Tassaddiq, M.S. Shabbir, Q. Din, H. Naaz, Discretization, Bifurcation, and Control for a Class of Predator-Prey Interactions, Fractal Fract. 6 (2022), 31. https://doi.org/10.3390/fractalfract6010031.
- [32] A.C.J. Luo, Regularity and Complexity in Dynamical Systems, Springer, New York, 2012. https://doi.org/10. 1007/978-1-4614-1524-4.
- [33] Y.A. Kuznetsov, Elements of Applied Bifurcation Theory, Springer, New York, 2004. https://doi.org/10.1007/ 978-1-4757-3978-7.
- [34] S. Wiggins, M. Golubitsky, Introduction to Applied Nonlinear Dynamical Systems and Chaos, Springer, New York, 2003. https://doi.org/10.1007/b97481.
- [35] S. Akhtar, R. Ahmed, M. Batool, N.A. Shah, J.D. Chung, Stability, Bifurcation and Chaos Control of a Discretized Leslie Prey-Predator Model, Chaos Solitons Fractals. 152 (2021), 111345. https://doi.org/10.1016/j.chaos. 2021.111345.

- [36] R. Ahmed, A. Ahmad, N. Ali, Stability Analysis and Neimark-Sacker Bifurcation of a Nonstandard Finite Difference Scheme for Lotka-Volterra Prey-Predator Model, Commun. Math. Biol. Neurosci. 2022 (2022), 61. https://doi. org/10.28919/cmbn/7534.
- [37] A. Al Khabyah, R. Ahmed, M.S. Akram, S. Akhtar, Stability, Bifurcation, and Chaos Control in a Discrete Predator-Prey Model With Strong Allee Effect, AIMS Math. 8 (2023), 8060–8081. https://doi.org/10.3934/math. 2023408.
- [38] M.S. Shabbir, Q. Din, R. Alabdan, A. Tassaddiq, K. Ahmad, Dynamical Complexity in a Class of Novel Discrete-Time Predator-Prey Interaction With Cannibalism, IEEE Access. 8 (2020), 100226–100240. https://doi.org/ 10.1109/access.2020.2995679.
- [39] A.A. Elsadany, Q. Din, S.M. Salman, Qualitative Properties and Bifurcations of Discrete-Time Bazykin-berezovskaya Predator-prey Model, Int. J. Biomath. 13 (2020), 2050040. https://doi.org/10.1142/s1793524520500400.