International Journal of Analysis and Applications



Determination of Sub-Diffusion Process and Source Term from Nonlocal Data: Applications to Microwave Radiations

Ilyas Khan^{1,2}, Arifa Samreen^{3,*}, Saqib Zia³, Ansar Javed³, Wei Sin Koh⁴

*Corresponding author: arifasamreen002@gmail.com

Abstract. We investigate two inverse problems (IPs) for the time fractional diffusion equation (TFDE) with an involution. The determination of a space varying source term, along with solution of a diffusion equation containing n fractional order derivatives in Caputo's sense from extra data at a specific time, constitutes the first IP. The second IP investigates the extracting of a time varying source term as well as the solution of the TFDE from non-local type extra condition. The second IP has applications to microwave radiations. The existence and uniqueness results for the solutions of both IPs are presented.

1. Introduction and formulation of the problems

We will define two IPs for the n terms TFDE with involution in domain $(z,t) \in \Omega$, where, $\Omega = \{(z,t) : -\pi < z < \pi, 0 < t < T\}$. The IPs are described for the following equation,

$${}^{C}D_{0|t}^{\alpha_{0}}\phi(z,t) + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}}\phi(z,t) - \phi_{zz}(z,t) + \epsilon\phi_{zz}(-z,t) = H(z,t), \tag{1.1}$$

Received: Mar. 23, 2025.

2020 Mathematics Subject Classification. 35R11.

Key words and phrases. fractional diffusion equation; inverse source problem; multinomial Mittag-Leffler function; Laplace transform; smart microgrids.

ISSN: 2291-8639

¹Department of Mathematics, Saveetha School of Engineering, SIMATS, Chennai, Tamil Nadu, India

²Department of Mathematics, College of Science Al-Zulfi, Majmaah University, Al-Majmaah 11952, Saudi Arabia

³Department of Mathematics, COMSATS University Islamabad, Park Road, Chak Shahzad Islamabad, Pakistan

⁴INTI International University, Persiaran Perdana BBN Putra Nilai, 71800 Nilai, Negeri Sembilan, Malaysia

where, ${}^CD^{\alpha_m}_{0|t}\phi(z,t)$ stands for Caputo fractional derivatives (CFDs) of order $0 < \alpha_m < 1$ and $\mu_m \ge 0$ where m = 0, 1, 2, ..., n are real constants, subject to the boundary conditions,

$$\phi(-\pi, t) = \phi(\pi, t), \quad t \in [0, T], \tag{1.2}$$

and initial condition,

$$\phi(z,0) = \psi(z), \quad z \in [-\pi, \pi],$$
 (1.3)

Physically, the mathematical method presented (1.1)-(1.4) describe the anomalous thermal diffusion cycle in a tight copper wire wrapped around a thin sheet of insulating material. It is assumed that the insulation layer is highly permeable. The temperature on one side has an influence on the diffusion process from the other. The standard diffusion equation is modified for this reason, and we add a third term with an involution. Such process leads to the consideration of an IP for an equation of n parameters TFDE and periodic boundary conditions with respect to a spatial variable, for more detail see [1]- [3]. The non-integer order derivatives play an important role in explaining physical process abnormalities. In this article, we consider multi-term fractional derivative in time because there are anomalies in diffusion/transport process. This non-standard behavior in diffusion/transport process can be described by many ways. One way is to describe these anomalies through many techniques already available in literature e.g. non-chaotic slicer map, continuous time random walk (CTRW), stochastic process etc [4]- [7]. Another way is to introduce the fractional derivative in corresponding equation. So our motivation to consider multi-terms fractional derivatives in time is that we will have more parameters to explain the anomalous behavior in diffusion or transport model. By considering more parameters we are able to fix these anomalies. If the order of fractional derivative lies between 0 and 1, then our model describe the sub-diffusion process.

We have another boundary condition by considering a process in such a way that temperature at one end for all the time t is proportional to the rate of fluctuation of the average speed value of temperature across the wire. Then,

$$\phi(-\pi, t) = \rho \left({}^{c}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{c}D_{0|t}^{\alpha_{m}} \right) \int_{-\pi}^{\pi} \phi(\zeta, t) d\zeta, \quad t \in [0, T].$$
(1.4)

Here, ρ denotes proportionality constant.

From nonlocal boundary condition (1.4), and also by using (1.1), we obtain

$$\phi(-\pi,t) = \rho \int_{-\pi}^{\pi} \left\{ \phi_{\zeta\zeta}(\zeta,t) - \epsilon \phi_{\zeta\zeta}(-\zeta,t) + H(\zeta,t) \right\} d\zeta, \quad t \in [0,T],$$

which implies that

$$\phi(-\pi,t) = \rho(1-\epsilon) \left[\phi_{\zeta}(\pi,t) - \phi_{\zeta}(-\pi,t) \right] + \rho \int_{-\pi}^{\pi} H(\zeta,t) d\zeta, \qquad t \in [0,T].$$

We define a new function

$$U(x,t) = \phi(x,t) - \rho \int_{-\pi}^{\pi} H(\zeta,t)d(\zeta),$$

Consequently, in terms of U(x,t), we have the following IP,

$${}^{C}D_{0|t}^{\alpha_{0}}U(z,t) + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}}U(z,t) - U_{zz}(z,t) + \epsilon U_{zz}(-z,t) = H(z,t), \tag{1.5}$$

and the boundary conditions become,

$$U_{z}(-\pi,t) - U_{z}(\pi,t) - bU(\pi,t) = 0, U(-\pi,t) - U(\pi,t) = 0.$$
(1.6)

Physically the background of TFDE with involution which is given in 1.5 is a variation of the diffusion equation that demonstrates non-local behavior and describes transport processes with extended memory. It is obtained by substituting the first-order time derivative in the standard diffusion equation with a fractional derivative of order β , where β belongs to the interval (0, 1) given in [8]. The physical interpretation of this equation involves the emergence of self-similarity combined with the long-time limit, which is connected to experimental observations [9]. The equation is employed to examine anomalous diffusion in intricate environments, taking into consideration the impact of a uniform external field on the dynamics of a particle [10]. Numerical analysis and simulations of the equation have been performed to comprehend sub-diffusive transport processes and verify the stability and convergence of numerical schemes [11]. The time fractional diffusion equation with involution is a mathematical model used to describe anomalous diffusion in various physical systems. It exhibits non-local behavior and is characterized by FDs. The physical interpretation of this equation is still an active area of research. Several papers have explored the physical background of the TFDE and its applications. Baeumer et al. discuss the derivation of physically meaningful boundary conditions for fractional diffusion equations and highlight the unsuitability of the CFD for modeling fractional diffusion [8]. Nadal et al. emphasize the non-local behavior of the fractional diffusion equation and its implications for physical understanding [13]. Bakalis and Zerbetto study the influence of a random moving particle in a complex environment using the TFDE and provide analytical solutions for various properties of the particle's motion [14].

The second condition in (1.6) refers to the inclusivity of the transmission intensity at the ends of the interval in a physical sense. While the first condition of (1.6) refers to the difference of flow rates along with opposite boundaries to the density value at the boundary in a physical sense. This Dirichlet type condition $U(\pi,t)=U(-\pi,t)=0$, was used instead of condition (1.6) in [15] alongside initial and final conditions

$$U(z,0) = \psi(z) - \rho \int_{-\pi}^{\pi} H(\zeta, t) d(\zeta), \quad z \in [-\pi, \pi],$$
(1.7)

$$U(z,T) = \Psi(z) - \rho \int_{-\pi}^{\pi} H(\zeta,t)d(\zeta), \quad z \in [-\pi,\pi], \tag{1.8}$$

and U(z,t) as a solution of Eq (1.5) satisfies (1.7)-(1.8), where $\psi(z)$ and $\Psi(z)$ are given smooth functions; such that $|\epsilon| < 1$; and $b = \frac{1}{\rho(\epsilon - 1)}$.

Physically IPs arise in almost all areas of science and technology, in modeling of problems motivated by various physical and social processes. Most of these models are governed by differential and integral equations. If all the necessary inputs in these models are known, then the solution can be computed and behavior of the physical system under various conditions can be predicted. In terms of differential problems, the necessary inputs include such information as initial or boundary data, coefficients and force term, also shape and size of the domain. If all these data are enough to describe the system adequately, then it is possible to use the mathematical model for studying the physical system [16].

The investigation into boundary-value problems that are not localized is prompted by the realization that, in numerous instances, a nonlocal condition is a more accurate approach to addressing physical problems in comparison to the conventional local conditions. The examination of IPs with nonlocal boundary conditions is explored, for instance, in the works of [17]- [19].

In [20]- [22], for space TFDE by using a novel modified quasi-reversibility regularization method, conjugate gradient method and non-stationary iterative Tikhonov regularization method have been used for simultaneous identification of source term and initial data. IPs for time fractional diffusion-wave equation have been considered in [23]- [25]. In the articles [26]- [27], IPs of stochastic as well as distributed order diffusion equation are discussed.

Prior research has looked into the well-posedness of direct and IPs for parabolic equations with involution [28]- [30]. The solvability of many IPs for parabolic equations was investigated in articles by Anikonov, Y. E., and Belov, Y. Y., Bubnov, B. A., Prilepko, A. I., and Kostin, A. B., Monakhov, V. N., Kozhanov, A. I., Kaliev, I. A., Sabitov, K. B., and several others [31]- [32]. In Ahmad et al, [15] there are good references to publications on these types of problem. The article [33]- [45] from literature are close to our article's theme. In these articles, various types of direct and inverse initial-boundary value problems for evolutionary equations are considered, including problems with nonlocal boundary conditions and problems for equations with fractional derivatives. Some of these articles use the concept of smart microgrids to solve inverse problems to understand, predict, and control their internal dynamics under uncertainty. For equation (1.1), we define two IPs.

Inverse Source Problem-I for TFDE (ISP-I): In the ISP-I the right hand side of the TFDE will be considered a space dependent source term. In ISP-I we define H(z,t):=h(z) along with U(z,t) which is the solution of the system (1.5)-(1.8) along with over-specified condition $\phi(z,T)=\Psi(z)$. The strong solution of the ISP-I means a couple of functions $\left\{U(z,t),h(z)\right\}$ that satisfies the system (1.5) – (1.8), such that $t^{\alpha_0+\alpha_m}U(z,t)\in C^{2,1}_{z,t}(\overline{\Omega}), t^{\alpha_0+\alpha_m}h(z)\in C([-\pi,\pi]),$ and $t^{\alpha_0+\alpha_m}\left({}^CD^{\alpha_0}_{0|t}+\sum_{m=1}^n\mu_m{}^CD^{\alpha_m}_{0|t}\right)U(z,t)\in C((0,T]),$ when $0<\alpha_m<1$ and μ_m is a real number. We proved that under certain regularity conditions on the given data the ISP-I has a unique strong solution.

Inverse Source Problem-II for TFDE (ISP-II): In the ISP-II, we want to recover time varying source term s(t) and U(z,t) for the system (1.5) - (1.8), with H(z,t) = s(t)h(z,t) for ISP-II. This structure of the source term arise in microwave heating process, in which the external energy is supplied to a target at a controlled level represented s(t) and h(z,t) is the local conversion rate of the microwave energy. The ISP-II does not give us unique solution, so for the uniqueness result of ISP-II we introduce an over-determination condition. Then, with the help of over-determination condition we are able to find the unique solution for the ISP-II. As, the over-determination condition is followed by

$$\int_{-\pi}^{\pi} U(z,t)dz = E(t). \tag{1.9}$$

The condition (1.9) is consider for the unique solution of the IP-II. ISP-II has a strong solution under some assumptions, that is, there exist a pair $\{U(z,t),s(t)\}$ with $t^{\alpha_0+\alpha_m}U(z,t)\in C^{2,1}_{z,t}(\overline{\Omega})$,

$$s(t) \in C([0,T])$$
, and $t^{\alpha_0 + \alpha_m} \left({}^{C}D_{0|t}^{\alpha_0} + \sum_{m=1}^n \mu_m {}^{C}D_{0|t}^{\alpha_m} \right) U(z,t) \in C((0,T])$.

The rest of the article is organized as fallows: Section 2 presents some basic concepts from Fractional Calculus (FC) and representation of multinomial Mittag-Leffler (ML) type functions and some its basic results are provided. The spectral problem, properties of eigenvalues and eigenfunctions for the IBVP (1.5)-(1.8) are described in Section 2.2. Section 3 provides the main results about existence and uniqueness of the solution of the ISP-I. In Section 4 we discuss the existence and uniqueness of the solution of ISP-II. Finally, in last section we conclude the paper and present some future perspectives.

2. Preliminaries

This section provides some preliminary results from FC which will be used in the forthcoming section. Indeed after providing definition of fractional order integrals and derivatives, we define the spectral problem (in space variable) of IBVP (1.5)-(1.8).

Let $g \in AC([a,b])$, i.e the space of absolutely continuous function then CFDs left and right sided of order $0 < \beta < 1$, are defined as

$${}^{C}D_{a_{+}|t}^{\beta}g(t) = \frac{1}{\Gamma(1-\beta)} \int_{a}^{t} (t-s)^{-\beta} \frac{d}{ds}g(s)ds,$$

$${}^{C}D_{b-|t}^{\beta}g(t) = \frac{-1}{\Gamma(1-\beta)} \int_{t}^{b} (s-t)^{-\beta} \frac{d}{ds}g(s)ds.$$

Where, for the convince of the reader, we provide the definition of absolutely continuous function Let $I \subset \mathbb{R}$. A function $g: I \to \mathbb{R}$ is absolutely continuous on I if for every $\epsilon > 0$, there is a positive number δ in a manner that whenever a sequence of disjoint sub-intervals (x_k, y_k) on interval I is in

a manner that $x_k < y_k \in I$ satisfies the following relation

$$\sum_{k} (y_k - x_k) < \delta$$

implies

$$\sum_{k} |g(y_k) - g(x_k)| < \epsilon.$$

2.1. Multinomial Mittag-Leffler Function. In this section, we are going to define multinomial Mittag-Leffler function and present some important results related to it which are used in the forthcoming sections.

The multinomial ML function, [48] for $\eta > 0$, $\xi_i > 0$, $z_i \in \mathbb{C}$, i = 0, 1, ..., m, $m \in \mathbb{N}$, is defined as where $(k; l_0, ..., l_m) = \frac{k!}{l_0! \times ... \times l_m!}$

Moreover, notice that

$$E_{(\xi_0,\xi_1,...,\xi_m),\eta}(z_0,z_1,...,z_m) = E_{(\xi_m,...,\xi_1,\xi_0),\eta}(z_m,...,z_1,z_0),$$
(2.1)

and

$$= \sum_{k=0}^{\infty} \sum_{l_0=0}^{k} \sum_{l_1=0}^{l_0} \dots \sum_{l_m=0}^{l_{m-1}} \frac{k!}{l_1! l_2! \dots (k-l_0-l_1-\dots-l_m)!} \times \frac{z_0^{l_0} z_1^{l_1} \dots z_m^{l_m}}{\Gamma(\eta + \xi_0 l_0 + \xi_1 l_1 + \dots + \xi_m l_m)}.$$

Remark 2.1. For $\xi_0 \neq 0$, and $z_0 \neq 0$, $z_1 = ... = z_m = 0$, $m \in \mathbb{N}$ the multinomial ML function become $E_{\xi_0,\eta}(z_0,0,0,...,0) = \sum_{k=0}^{\infty} \frac{z_0^k}{\Gamma(\eta+\xi_0 k)} := E_{\xi_0,\eta}(z_0)$, where $E_{\xi_0,\eta}(z_0)$ is the ML function of two parameters.

For n = 2 the multinomial ML written in series form

$$E_{(\xi_{0},\xi_{1}),\eta}(z_{0},z_{1}) = \sum_{k=0}^{\infty} \sum_{\substack{l_{0}+l_{1}=k\\l_{0}\geq 0, l_{1}\geq 0}} \frac{k!}{l_{0}! \, l_{1}!} \frac{z_{0}^{l_{0}} z_{1}^{l_{1}}}{\Gamma(\eta+\xi_{0}l_{0}+\xi_{1}l_{1})}$$

$$= \sum_{k=0}^{\infty} \sum_{l_{0}=0}^{k} \frac{k!}{l_{0}! \, (k-l_{0})!} \frac{z_{0}^{l_{0}} z_{1}^{k-l_{0}}}{\Gamma(\eta+\xi_{0}l_{0}+\xi_{1}(k-l_{0}))}.$$
(2.2)

For $z_1 = 0$, $\eta = 1$ the multinomial ML function takes form $E_{\xi_0,1}(z_0,0) = E_{\xi_0,1}(z_0)$, i.e., the one parameter ML function. Following Lemma 3.2 of, [49] and from 2.1

$$E_{(\xi_0,\xi_0-\xi_1,...,\xi_0-\xi_m),\eta}(z_0,z_1,...,z_m) = E_{(\xi_0-\xi_m,...,\xi_0-\xi_1,\xi_0),\eta}(z_m,...,z_1,z_0).$$

we can have the following Lemmata

Lemma 2.1. [50] For $0 < \eta < 1$ and $0 < \xi_m < ... < \xi_0 < 1$ be given. Assume that $\xi_0 \pi/2 < \mu < \xi_0 \pi$, $\mu \le |arg(z_m)| \le \pi$ and $z_i > 0$, i = 0, 1, ..., m. Then, there exists a constant depending only on $\mu, \xi_i, i = 0, 1, ..., m$ such that

$$\left| E_{(\xi_0 - \xi_m, ..., \xi_0 - \xi_1, \xi_0), \eta}(z_m, ..., z_1, z_0) \right| \le \frac{C_1}{1 + |z_m|}.$$

We use the following notations throughout the paper

$$\Xi := (\xi_0 - \xi_1, \xi_0 - \xi_2..., \xi_0 - \xi_n, \xi_0)$$

and

$$\mathcal{E}_{\Xi,\eta}(t;m_1,m_2,...,m_n) := t^{\eta-1} E_{\Xi,\eta}(-m_1 t^{\xi_0-\xi_1},-m_2 t^{\xi_0-\xi_2},...,-m_n t^{\xi_0}),$$

where m_i are positive real constants.

Lemma 2.2. [50] For η , $m_i > 0$, $\xi_i > 0$ and $0 < \xi_m < ... < \xi_0 < 1$ for i = 0, 1, ..., n, $n \in \mathbb{N}$ the Laplace transform of the multinomial ML function is given by

$$\mathcal{L}\left\{\mathcal{E}_{\Xi,\eta}\left(t;m_{1},m_{2},...m_{n-1},m_{n}\right);s\right\} = \frac{s^{\xi_{0}-\eta}}{\left(s^{\xi_{0}}+m_{1}s^{\xi_{1}}+...+m_{n-1}s^{\xi_{n}}+m_{n}\right)},\tag{2.3}$$

 $\text{if} \left| m_1 s^{\xi_0 - \xi_1} + m_2 s^{\xi_0 - \xi_2} + \ldots + m_{n-1} s^{\xi_0 - \xi_n} + m_n s^{\xi_0} \right| < 1.$

Lemma 2.3. [50] For $g \in C^1([a,b])$ and $m_i > 0$, $\lambda_i > 0$, for i = 0, 1, ..., n, we have

$$|g(t) * \mathcal{E}_{\Xi,\xi_0+1}(t; m_0, m_1, ..., m_n)| \le \frac{C_1}{m_n} ||g||_{C^1([0,T])},$$

where $\|.\|$ represents a norm given by $\|g\|_{C^1([0,T])} := \max_{0 \le t \le T} |g(t)| + \max_{0 \le t \le T} |g'(t)|$ and $C^1([0,T])$ denotes the space of a continuous first-order differentiable function.

2.2. **The spectral problem.** A spectral problem emerges when the Fourier method is used to address the problem (1.5)-(1.8). The spectral problem for the operator £ given by the following equation

$$\pounds X(z) \equiv -X''(z) + \epsilon X''(-z) = \lambda X(z), \quad -\pi < z < \pi,$$
 (2.4)

where, λ is a spectral parameter, and the boundary conditions are

$$X'(-\pi) - X'(\pi) = bX(\pi), X(-\pi) = X(\pi).$$
(2.5)

The equation (2.4) and equation (2.5) has two series of the eigenvalues

$$\lambda_{k1} = (1 + \epsilon)k^2, \quad k \in \mathbb{N},$$

and

$$\lambda_{k2} = (1 - \epsilon)(k + \sigma_k)^2, \quad \sigma_k = \frac{b}{k+1}O(1) > 0, \ k \in \mathbb{N}_0 \equiv \mathbb{N} \cup \{0\},$$

with associating normalized eigenfunctions given by

$$X_{k1} = \frac{1}{\sqrt{\pi}} \sin(kx), \ k \in \mathbb{N}, \ X_{k2} = w_k \cos((k + \sigma_k)x), \ k \in \mathbb{N}_0.$$
 (2.6)

Here, w_k denote the normalization coefficient:

$$w_k^{-2} = \|\cos((k+\sigma_k)x)\|^2 = \pi + \frac{b^2}{(k+\sigma_k)[b^2 + (k+\sigma_k)^2\pi^2]}.$$

3. The main results

In this section, we will provide our main results about both IPs defined for TFDE with involution. Indeed, this section has been divided into two subsections corresponding to the results of ISP-I and ISP-II.

3.1. **Inverse source problem-I (ISP-I).** In ISP-I, we consider H(z,t) = h(z). By eigenfunction expansion method the solution of the ISP-I can be written as

$$U(z,t) = U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty} \left[\sum_{i=1}^{2} U_{ki}(t)X_{ki}(z) \right], \tag{3.1}$$

where,
$$U_{02}(t) = \langle U(.,t), X_{02}(.) \rangle$$
, $X_{02}(z) = \left\{ \pi + \frac{1}{b(1+\pi^2)} \right\} \cos(bz)$ and $U_{ki}(t) = \langle U(.,t), X_{ki}(.) \rangle$.

Expression of h(z) is given by

$$h(z) = h_{02}X_{02}(z) + \sum_{k=1}^{\infty} \left[\sum_{i=1}^{2} h_{ki}X_{ki}(z) \right],$$
 (3.2)

where, $h_{02} = \langle h(.), X_{02}(.) \rangle$ and $h_{ki} = \langle h(.), X_{ki}(.) \rangle$.

Consider

$$U_{ki}(t) = \langle U(.,t), X_{ki}(.) \rangle, \qquad i = 1, 2.$$
 (3.3)

We apply the operator $\left({}^{C}D_{0|t}^{\alpha_0} + \sum_{m=1}^{n} \mu_m{}^{C}D_{0|t}^{\alpha_m}\right)$ to (3.3). Then, using Equation (1.1) and integration by parts, we obtain a n term TFDE

$$\left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}}}\right) U_{ki}(t) + \lambda_{ki} U_{ki}(t) = h_{ki}, \quad i = 1, 2,$$
(3.4)

$$U_{ki}(0) = \psi_{ki} - \delta_{ki}, \quad i = 1, 2,$$
 (3.5)

$$U_{ki}(T) = \Psi_{ki} - \delta_{ki}, \quad i = 1, 2.$$
 (3.6)

In above equations, we use notations

$$\psi_{ki} = \int_{-\pi}^{\pi} \psi(x) X_{ki}(z) d(z), \qquad \Psi_{ki} = \int_{-\pi}^{\pi} \Psi(z) X_{ki}(z) d(z),$$

$$\delta_{ki} = \rho \int_{-\pi}^{\pi} h(\zeta) d(\zeta) \int_{-\pi}^{\pi} X_{ki}(z) d(z).$$

After applying Laplace transformation on equation (3.4), we obtain

$$U_{ki}(s) = \frac{s^{\alpha_0 - 1}}{s^{\alpha_0} + \sum_{m=1}^{n} \mu_m s^{\alpha_m} + \lambda_{ki}} U_{ki}(0) + \sum_{m=1}^{n} \mu_m \frac{s^{\alpha_m - 1}}{s^{\alpha_0} + \sum_{m=1}^{n} \mu_m s^{\alpha_m} + \lambda_{ki}} U_{ki}(0) + h_{ki} \frac{1}{s \left(s^{\alpha_0} + \sum_{m=1}^{n} \mu_m s^{\alpha_m} + \lambda_{ki}\right)}.$$
(3.7)

Now apply inverse Laplace transformation on equation (3.7) and by virtue of Lemma 2.2, we have

$$U_{ki}(t) = \mathcal{E}_{\Xi,1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n})U_{ki}(0) + \sum_{m=1}^{n} \mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n})U_{ki}(0) + h_{ki}\mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}),$$
(3.8)

where, $U_{ki}(0)$ and h_{ki} are unknowns. By using initial condition we obtain

$$U_{ki}(t) = \psi_{ki} \bigg\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \\
+ \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \bigg\} \\
- \delta_{ki} \bigg\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \\
+ \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \bigg\} \\
+ h_{ki} \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}). \tag{3.9}$$

By, using the value of (3.6), we obtain

$$h_{ki} = \frac{1}{\mathcal{E}_{\Xi,\alpha_0+1}(T;\lambda_{ki},\mu_1,\mu_2,...,\mu_n)} \times \left[(\Psi_{ki} - \delta_{ki}) - \psi_{ki} \Big\{ \mathcal{E}_{\Xi,1}(T;\lambda_{ki},\mu_1,\mu_2,...,\mu_n) + \sum_{m=1}^{n} \mu_m \mathcal{E}_{\Xi,\alpha_0-\alpha_m+1}(T;\lambda_{ki},\mu_1,\mu_2,...,\mu_n) \Big\} - \delta_{ki} \Big\{ \mathcal{E}_{\Xi,1}(T;\lambda_{ki},\mu_1,\mu_2,...,\mu_n) \right]$$

$$+\sum_{m=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n})\bigg\}\bigg].$$
(3.10)

By putting values of (3.9) and (3.10) in (3.1) and (3.2), we obtain

$$U(z,t) = U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty} \left[\sum_{i=1}^{2} \left(\psi_{ki} \left\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right. \right. \right. \\ + \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} \\ - \delta_{ki} \left\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right. \\ + \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} \\ + h_{ki} \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \left. \right) X_{ki}(z) \right].$$

$$h(z) = h_{02} X_{02}(z) + \sum_{k=1}^{\infty} \left[\sum_{i=1}^{2} \frac{1}{\mathcal{E}_{\Xi,\alpha_{0}+1}(T; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n})} \times \left[(\Psi_{ki} - \delta_{ki}) - \psi_{ki} \right\} \mathcal{E}_{\Xi,1}(T; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right]$$

$$h(z) = h_{02}X_{02}(z) + \sum_{k=1} \left[\sum_{i=1}^{n} \frac{1}{\mathcal{E}_{\Xi,\alpha_{0}+1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n})} \times \right]$$

$$\left[(\Psi_{ki} - \delta_{ki}) - \psi_{ki} \left\{ \mathcal{E}_{\Xi,1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}) + \sum_{m=1}^{n} \mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}) \right\}$$

$$-\delta_{ki} \left\{ \mathcal{E}_{\Xi,1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}) + \sum_{m=1}^{n} \mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(T;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}) \right\} X_{ki}(z) \right].$$

$$(3.12)$$

Hence, the couple of functions $\{U(z,t),h(z)\}$ is required solution for the ISP-I.

3.2. **Existence of the Solution of the ISP-I.** In this subsection, we will prove a classical solution for the ISP-I, from the following theorem.

Theorem 3.1. If $\psi(z)$ and $\varphi(z) \in C^4([-\pi, \pi])$ and $\psi(z)$, $\varphi(z)$ and $\psi''(z)$, $\varphi''(z)$ satisfy the boundary conditions (1.6), then the ISP-I has a unique classical solution.

From the analysis of (3.10), it's easy to see that the solution (3.1) of the ISP-I, will form a convergent series iff

$$\lim_{k \to \infty} \delta_{ki} \left\{ \mathcal{E}_{\Xi,1}(T; \lambda_{ki}, \mu_1, \mu_2, ..., \mu_n) + \sum_{m=1}^{n} \mu_m \mathcal{E}_{\Xi,\alpha_0 - \alpha_m + 1}(T; \lambda_{ki}, \mu_1, \mu_2, ..., \mu_n) \right\} = 0. \qquad i = 1, 2.$$
(3.13)

As, we know that

$$\lim_{k\to\infty}\delta_{ki}=\lim_{k\to\infty}\left(\rho\int_{-\pi}^{\pi}h(\zeta)d\zeta\right)\int_{-\pi}^{\pi}X_{ki}(z)dz,\quad i=1,2.$$

which implies that

$$\lim_{k \to \infty} \delta_{ki} = \lim_{k \to \infty} \frac{d}{\lambda_{ki}} \left(\rho \int_{-\pi}^{\pi} h(\zeta) d\zeta \right) X_{ki}(\pi) dz = 0, \ i = 1, 2, 3, ..., m.$$
 (3.14)

Now, we shall show of U(z,t) given by (3.11) represents a continuous function. By using Lemma 2.1, we obtained

$$T^{\alpha_{0}+\alpha_{m}}|h(z)| \leq \frac{|\lambda_{02}|}{c_{1}}T^{\alpha_{0}+\alpha_{m}}|\varphi_{02}| - |\psi_{02}| \left\{ T^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m}T^{\alpha_{0}} \right\}$$

$$+ \sum_{k=1}^{\infty} \sum_{i=1}^{2} \left[\frac{|\lambda_{ki}|}{c_{1}} T^{\alpha_{0}+\alpha_{m}} |\varphi_{ki}| \right.$$

$$- |\psi_{ki}| \left\{ T^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m}T^{\alpha_{0}} \right\} \right],$$

implies that

Consider

$$T^{\alpha_{0}+\alpha_{m}}|h(z)| \leq \frac{|\lambda_{02}|}{c_{1}}T^{\alpha_{0}+\alpha_{m}}|\varphi_{02}| - |\psi_{02}| \left\{ T^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m}T^{\alpha_{0}} \right\}$$

$$+ \sum_{k=1}^{\infty} \left[T^{\alpha_{0}+\alpha_{m}}|\varphi^{4}(z)| \left\{ \frac{1+\epsilon}{k^{2}} + \frac{1-\epsilon}{(k+\sigma_{k})^{2}} \right\} \right]$$

$$- \sum_{i=1}^{2} |\psi_{ki}| \left\{ T^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m}T^{\alpha_{0}} \right\} \right].$$

Hence, $T^{\alpha_0 + \alpha_m} h(z)$ is uniformly convergent. Now we are going to prove that the series related to the operator $\left({}^C D_{0|t}^{\alpha_0} + \sum_{m=1}^n \mu_m {}^C D_{0|t}^{\alpha_m}\right) U(z,t)$ also uniformly convergent.

$$\left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}}}\right)U(z,t) = \left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}}}\right)U_{02}(t)X_{02}(z)$$

$$+\left({}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}}\right)$$
$$\sum_{k=1}^{\infty} \sum_{i=1}^{2} \left[U_{ki}(t) X_{ki}(z) \right], \tag{3.15}$$

by using Lemma 15.2 (page 278, see [47])

$$\begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U(z,t) = \begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U_{02}(t)X_{02}(z) \\
+ \sum_{k=1}^{\infty} \sum_{i=1}^{2} \begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} \\
U_{ki}(t)X_{ki}(z), \\
= h(z) - \left[\lambda_{02}U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty} \sum_{i=1}^{2} \lambda_{ki} \{U_{ki}(t)X_{ki}(z)\}\right].$$

First, we will prove that the series $\left[\lambda_{02}U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty}\sum_{i=1}^{2}\lambda_{ki}\left\{U_{ki}(t)X_{ki}(z)\right\}\right]$ is uniformly convergent, by using Lemma 2.1, we obtain

$$t^{\alpha_{0}+\alpha_{m}} \left| \sum_{k=1}^{\infty} \sum_{i=1}^{2} \lambda_{ki} U_{ki}(t) X_{ki}(z) \right| \leq \sum_{k=1}^{\infty} c_{1} |\psi^{4}(z)| \left(\frac{1+\epsilon}{k^{2}} + \frac{1-\epsilon}{(k+\sigma_{k})^{2}} \right) \left\{ t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right\} + \sum_{i=1}^{2} t^{\alpha_{0}+\alpha_{m}} |h_{ki}|,$$

and

$$t^{\alpha_0+\alpha_m} |\lambda_{02} U_{02}(t) X_{02}(z)| \le c_1 |\psi_{02}| \left\{ t^{\alpha_m} + \sum_{m=1}^n \mu t^{\alpha_0} \right\} + t^{\alpha_0+\alpha_m} |h_{02}|.$$

By virtue of Weierstrass M-test the series $t^{\alpha_0 + \alpha_m} \left[\lambda_{02} U_{02}(t) X_{02}(z) + \sum_{k=1}^{\infty} \sum_{i=1}^{2} \frac{1}{i-1} \right]$

 $\lambda_{ki}\{U_{ki}(t)X_{ki}(z)\}$ represents a continuous function. Furthermore, the uniform convergence of $T^{\alpha_0+\alpha_m}h(z)$ have been proved already. First, we will prove the uniqueness result for the ISP-I.

Theorem 3.2. Consider a couple of functions $\{U_1(z,t),h_1(z)\}$ and $\{U_2(z,t),h_2(z)\}$ be the two classical solution sets for the ISP-I and $z_0 \in (-\pi,\pi)$, such that $U_1(z_0,t) = U_2(z_0,t)$, then $U_1(z,t) = U_2(z,t)$ implies $h_1(z) = h_2(z)$ for all $z \in (-\pi,\pi)$.

Proof. The proof of theorem 3.2 can be proved by following the same strategy given in [51].

4. ISP-II

In this section, we will introduce a pair of function $\{U(z,t),s(t)h(z,t)\}$ for the system (1.5)-(1.8), for the ISP-II. Whenever, the over-determination condition (1.9) is given. Furthermore, we will be proved s(t), $t^{\alpha_0+\alpha_m}U(z,t)$ and $t^{\alpha_0+\alpha_m}\binom{CD_{0|t}^{\alpha_0}}{t} + \sum_{m=1}^n \mu_m{}^CD_{0|t}^{\alpha_m}\binom{DD_{0|t}^{\alpha_m}}{t} U(z,t)$ represents a continuous functions.

4.1. **Solution for the ISP-II.** By using eigen-function expansion method the solution for the ISP-II can be written by

$$U_{ki}(t) = \langle U(.,t), X_{ki}(.) \rangle$$
 $i = 1, 2.$ (4.1)

where $U_{ki}(t)$ satisfy the following fractional differential equation

$$\left({^{C}D_{0|t}^{\alpha_0} + \sum_{m=1}^{n} \mu_m {^{C}D_{0|t}^{\alpha_m}}}\right) U_{ki}(t) + \lambda_{ki} U_{ki}(t) = s(t)h_{ki}(t), \ i = 1, 2.$$
(4.2)

By following the same strategy which we used for ISP-I, we find $U_{ki}(t)$, then we write the series solution for ISP-II

$$U_{ki}(t) = \psi_{ki} \Big\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \\
+ \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \Big\} \\
- \delta_{ki} \Big\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \\
+ \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \Big\} \\
+ \Big\{ s(t)h_{ki}(t) \Big\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}).$$

$$U(z,t) = \Big[\Big\{ \psi_{02} \Big\} \mathcal{E}_{\Xi,1}(t; \lambda_{02}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \Big\}$$

$$(4.3)$$

$$U(z,t) = \left[\left(\psi_{02} \left\{ \mathcal{E}_{\Xi,1}(t; \lambda_{02}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right. \right. + \left. \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{02}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} \right. \\ \left. + \left\{ s(t)h_{02}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{02}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} X_{02}(z) \right] \\ \left. + \sum_{k=1}^{\infty} \sum_{i=1}^{2} \left[\left(\psi_{ki} \left\{ \mathcal{E}_{\Xi,1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) + \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) + \left\{ s(t)h_{ki}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} X_{ki}(z) \right],$$

$$\left. + \left\{ s(t)h_{ki}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t; \lambda_{ki}, \mu_{1}, \mu_{2}, ..., \mu_{n}) \right\} X_{ki}(z) \right],$$

$$\left. (4.4)$$

where s(t) is still to be determined.

4.2. **Existence of the Solution for the ISP-II.** Now, we proved the existence of solution for ISP-II in the given domain $\overline{\Omega}$, under the assumptions of following theorem:

Theorem 4.1. The following conditions should be satisfied for the existence of solution for the ISP-II.

- (i) $\psi \in C^2([-\pi, \pi])$ be such that $\psi(-\pi) = 0 = \psi(\pi)$.
- (ii) $h(.,t) \in C^2([-\pi,\pi])$ be such that $h(-\pi,t) = 0 = h(\pi,t)$. Moreover,

$$\int_{-\pi}^{\pi} h(z,t)dz \neq 0$$

and

$$\left(\int_{-\pi}^{\pi} h(z,t)dz\right)^{-1} \le M_1, \quad \forall \quad 0 \le t \le T.$$

Where, the constant $M_1 > 0$.

(iii) $E \in AC([0,T])$ and E(t) satisfies the condition

$$\int_{-\pi}^{\pi} \psi_m(z) dz = E(0). \qquad m = 0, 1, 2, ..., n.$$

Then, there exist a unique classical solution for the ISP-II.

Proof: To find s(t), we will use extra condition (1.9) and apply the operator $\left({}^{C}D_{0|t}^{\alpha_0} + \sum_{m=1}^{n} \mu_m {}^{C}D_{0|t}^{\alpha_m}\right)$ to (1.9), then we have

$$\int_{-\pi}^{\pi} \left({^{C}D_{0|t}^{\alpha_{0}}U(z,t) + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}U(z,t)} \right) dz$$

$$= \left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}} \right) E(t),$$
(4.5)

from equation (1.5), we obtain

$$\int_{-\pi}^{\pi} \left(U_{zz}(z,t) - \epsilon U_{zz}(z,t) \right) dz + s(t) \int_{-\pi}^{\pi} h(z,t) dz$$

$$= \left({}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m} {}^{C}D_{0|t}^{\alpha_{m}} \right) E(t), \tag{4.6}$$

this results in the following: Integral equation of the Volterra kind

$$s(t) = \frac{\left({}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}}\right) E(t) - \int_{-\pi}^{\pi} \left(U_{zz}(z,t) - \epsilon U_{zz}(-z,t)\right) dz}{\int_{-\pi}^{\pi} h(z,t) dz},$$
(4.7)

first, we solve it

$$\int_{-\pi}^{\pi} \left(U_{zz}(z,t) - \epsilon U_{zz}(-z,t) \right) dz, \tag{4.8}$$

after taking integration of (4.8) and then using boundary conditions, we obtain $d(1+\epsilon)U(\pi,t)$. Similarly, we have another way to solve equation (4.8), by using Fourier method,

$$T(t)\int_{-\pi}^{\pi} \left\{X^{\prime\prime}(z) - X^{\prime\prime}(-z)\right\} dz,$$

by taking integration and boundary conditions (2.5), we obtain

$$d(1+\epsilon)T(t)X(\pi) \Rightarrow d(1+\epsilon)U(\pi,t).$$

which, is same in both cases.

Then, we have

$$d(1+\epsilon)U(\pi,t)$$

$$= d(1+\epsilon) \left[\left(\psi_{02} \left\{ \mathcal{E}_{\Xi,1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n}) + \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n}) \right\} + \left\{ s(t)h_{02}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n}) \right\} X_{02}(\pi) \right]$$

$$+ \sum_{k=1}^{\infty} \left[\left(\psi_{k2} \left\{ \mathcal{E}_{\Xi,1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n}) + \sum_{m=1}^{n} \mu_{n} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n}) + \left\{ s(t)h_{k2}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n}) \right\} X_{k2}(\pi) \right],$$

$$(4.9)$$

which implies that

$$s(t) = \left(\int_{-\pi}^{\pi} g(z,t)dz\right)^{-1} \left[\left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{^{C}D_{0|t}^{\alpha_{m}}}}\right)E(t)\right]$$

(4.11)

$$+d(1+\epsilon)\left(\psi_{02}\left\{\mathcal{E}_{\Xi,1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n})\right.\right.$$

$$+\sum_{m=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n})\right\}$$

$$+d(1+\epsilon)\sum_{k=1}^{\infty}\left[\left(\psi_{k2}\left\{\mathcal{E}_{\Xi,1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})\right.\right.\right]$$

$$+\sum_{m=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})\right]$$

$$+\left\{h_{k2}(t)\right\}*\mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})\right\}$$

$$+\left\{h_{k2}(t)\right\}*\mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})\right\}$$

$$+\sum_{m=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{02},\mu_{1},\mu_{2},...,\mu_{n})$$

$$+d(1+\epsilon)\sum_{k=1}^{\infty}\left\{\psi_{k2}\left\{\mathcal{E}_{\Xi,1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})\right\}\right\}$$

$$+\sum_{k=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})$$

$$+\sum_{k=1}^{n}\mu_{m}\mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{k2},\mu_{1},\mu_{2},...,\mu_{n})$$

and

$$k(t,\tau) = d(1+\epsilon) \Big(\Big\{ h_{02}(\tau) \Big\} * \mathcal{E}_{\Xi,\alpha_0+1}(t-\tau;\lambda_{02},\mu_1,\mu_2,...,\mu_n) \Big) X_{02}(\pi)$$

$$+ d(1+\epsilon) \sum_{k=1}^{\infty} \Big(\Big\{ h_{k2}(\tau) \Big\} * \mathcal{E}_{\Xi,\alpha_0+1}(t-\tau;\lambda_{k2},\mu_1,\mu_2,...,\mu_n) \Big)$$

$$\cos((k+\delta_k)\pi).$$
(4.12)

Then equation (4.10) can be written as,

$$s(t) = \left(\int_{-\pi}^{\pi} g(z,t) dx \right)^{-1} \left[\left({^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m} {^{C}D_{0|t}^{\alpha_{m}}}} \right) E(t) + \mathcal{F}(t) \right]$$
$$+ \int_{0}^{t} s(\tau)k(t,\tau) d\tau d\tau.$$

 $\cos((k+\delta_k)\pi)$,

Now, we consider a space of continuous functions C([0,T]), with the Chebyshev norm. Then, we are going to define a map $M: C([0,T]) \to C([0,T])$ by

$$M(s(t)) := s(t),$$

where, s(t) is given by (4.7).

$$t^{\alpha_{0}+\alpha_{m}}|F(t)| \leq d(1+\epsilon)\frac{c_{1}|\psi_{02}|}{|\lambda_{02}|} \left[\left\{ t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right\} \right] + d(1+\epsilon)\frac{c_{1}|\psi_{k2}|}{|\lambda_{k2}|} \left[\left\{ t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right\} \right]. \tag{4.13}$$

The uniform convergence of the series involves in (4.13) is ensured by using Lemma 2.1 and continuity of $|\psi_{k2}|$. Hence, by Weierstrass M-test F(t) presents a continuous function. Now, we shall show that equation (4.12), represents a continuous function.

$$|k(t,\tau)| \le d(1+\epsilon) \left[\frac{c_1}{|\lambda_{02}|} ||h_{02}|| + \frac{c_1}{|\lambda_{k2}|} ||h_{k2}|| \right].$$
 (4.14)

Then, the uniform convergence of the series involves in (4.14) is ensured by using Lemma 2.3 and continuity of g(z,t). Hence, by Weierstrass M-test $k(t,\tau)$, represents a continuous function such that

$$|K(t,\tau)| \leq M_2. \tag{4.15}$$

Hence, the mapping is well defined. Now, we will show that M(s(t)) := s(t), is a contraction mapping under the assumption $T < \frac{1}{M_1 M_2}$. For this we consider

$$|\mathbb{M}(s(t)) - \mathbb{M}(c(t))| \le \left(\int_{-\pi}^{\pi} h(z,t)dz\right)^{-1} \int_{0}^{t} k(t,\tau)|s(\tau) - c(\tau)|d\tau,$$

By assumptions of the Theorem 4.1, we obtain

$$\max_{0 \le t \le T} |M(s(t)) - M(c(t))| \le M_1 M_2 T \max_{0 \le t \le T} |s(\tau) - c(\tau)|.$$

which implies

$$||M(s(t)) - M(c(t))||_{C([0,T])} \le M_1 M_2 T ||s - c||_{C([0,T])}.$$

Where, M_1 and M_2 are positive constants independents of n and M_2 is given by (4.15) which implies that the mapping $M(\cdot)$ is contraction map. Hence, by Banach fixed point theorem unique existence of s(t) is proved.

By using Lemma 2.1 and Lemma 2.3 in equation (4.3), we have the following relation

$$t^{\alpha_{0}+\alpha_{m}}|U_{ki}(t)| \leq \frac{c_{1}}{|\lambda_{ki}|} \left\{ |\psi_{ki}| \left(t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right) + t^{\alpha_{0}+\alpha_{m}} ||sh_{ki}||_{C([0,T])} \right\}$$

$$(4.16)$$

Since $s \in C([0,T])$, $||s||_{C([0,T])} < M$ where M > 0, and by using Lemma 2.3

$$t^{\alpha_{0}+\alpha_{m}}|U_{ki}(t)| \leq \frac{c_{1}}{|\lambda_{ki}|} \left\{ |\psi^{4}(z)| \left(\frac{1+\epsilon}{k^{2}} + \frac{1-\epsilon}{(k+\sigma_{k})^{2}} \right) \left(t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right) + t^{\alpha_{0}+\alpha_{m}} M ||h_{ki}||_{C([0,T])} \right\}$$

Next, we will prove the continuity of solution U(z,t), which is given by (4.4) and each term of (1.5), that is $U_{zz}(z,t)$, $U_{zz}(-z,t)$.

By using Lemma 2.1 and Lemma 2.3 in equation (4.4), we obtain

$$t^{\alpha_{0}+\alpha_{m}}|U(z,t)| \leq \frac{c_{1}}{\lambda_{02}} \left[|\psi_{02}| \left\{ t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right\} + t^{\alpha_{0}+\alpha_{m}} M ||h_{02}||_{C([0,T])} \right]$$

$$+ \sum_{k=1}^{\infty} \sum_{i=1}^{2} \frac{c_{1}}{|\lambda_{ki}|} |\psi^{4}(z)| \left(\frac{1+\epsilon}{k^{2}} + \frac{1-\epsilon}{(k+\sigma_{k})^{2}} \right)$$

$$\left(t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right) + M ||h_{ki}||_{C([0,T])} t^{\alpha_{0}+\alpha_{m}} \right].$$

By Weierstrass M-test the above series represent a continuous function.

Now, we are going to prove that the series related to the operator $t^{\alpha_0 + \alpha_m} \left({}^C D_{0|t}^{\alpha_0} + \sum_{m=1}^n \mu_m {}^C D_{0|t}^{\alpha_m} \right) U(z,t)$ also uniformly convergent, for this we consider the following relation

$$\begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U(z,t) = \begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U_{02}(t) X_{02}(z) \\
+ \begin{pmatrix}
{}^{C}D_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} \\
\sum_{k=1}^{\infty} \sum_{i=1}^{2} \left[U_{ki}(t) X_{ki}(z) \right], \tag{4.17}$$

by using Lemma 15.2 (page 278, see [47]), we have

$$\begin{pmatrix}
CD_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U(z,t) = \begin{pmatrix}
CD_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} U_{02}(t)X_{02}(z) \\
+ \sum_{k=1}^{\infty} \sum_{i=1}^{2} \begin{pmatrix}
CD_{0|t}^{\alpha_{0}} + \sum_{m=1}^{n} \mu_{m}^{C}D_{0|t}^{\alpha_{m}} \end{pmatrix} \\
\left[U_{ki}(t)X_{ki}(z)\right] \\
= s(t)h(z,t) - \left[\lambda_{02}U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty} \sum_{i=1}^{2} \lambda_{ki} \left\{U_{ki}(t)X_{ki}(z)\right\}\right],$$

we already prove the local well-posedness of s(t), we will prove that the series $\left[\lambda_{02}U_{02}(t)X_{02}(z) + \sum_{k=1}^{\infty}\sum_{i=1}^{2}\lambda_{ki}\left\{U_{ki}(t)X_{ki}(z)\right\}\right]$ is uniformly convergent, by using Lemma 2.1 and 2.3, we obtain

$$t^{\alpha_{0}+\alpha_{m}} \left| \sum_{k=1}^{\infty} \sum_{i=1}^{2} \lambda_{ki} U_{ki}(t) X_{ki}(z) \right| \leq \sum_{k=1}^{\infty} c_{1} |\psi^{4}(z)| \left(\frac{1+\epsilon}{k^{2}} + \frac{1-\epsilon}{(k+\sigma_{k})^{2}} \right)$$

$$\left\{ t^{\alpha_{m}} + \sum_{m=1}^{n} \mu_{m} t^{\alpha_{0}} \right\} + \sum_{i=1}^{2} t^{\alpha_{0}+\alpha_{m}}$$

$$M \|h_{ki}\|_{C([0,T])},$$

and

$$t^{\alpha_0 + \alpha_m} |\lambda_{02} U_{02}(t) X_{02}(z)| \leq c_1 |\psi_{02}| \left\{ t^{\alpha_m} + \sum_{m=1}^n \mu_m t^{\alpha_0} \right\} + t^{\alpha_0 + \alpha_m} M ||h_{02}||_{C([0,T])}.$$

Hence, the above series converge uniformly. Similarly we proved $U_{zz}(z,t)$, $U_{zz}(-z,t)$ represent a continuous function.

4.3. **Uniqueness of the Solution of ISP-II.** In this Subsection, we will consider the uniqueness of the ISP-II solution.

Theorem 4.2. *Under the assumptions of Theorem 4.1, solution of the ISP-II is unique.*

Proof: With the help of Banach fixed point theorem, we prove the uniqueness of the source term s(t). Now, we will present uniqueness of U(z,t).

Let U(z,t) and V(z,t) be the two regular solution sets of the ISP-II, and let $\widetilde{U}(z,t) = U(z,t) - V(z,t)$. Then, the function $\widetilde{U}(z,t)$ satisfy the following equation

$${}^{C}D_{0|t}^{\alpha_{0}}\widetilde{U}(z,t) + \sum_{m=1}^{n} \mu_{m}{}^{C}D_{0|t}^{\alpha_{m}}\widetilde{U}(z,t) - \widetilde{U}_{zz}(z,t) + \widetilde{\epsilon}\widetilde{U}_{zz}(-z,t) = 0,$$

with the boundary conditions (1.6) and the initial condition

$$\widetilde{U}(z,0) = 0, (4.18)$$

Consider the functions

$$\widetilde{U}_{ki}(t) = \langle \widetilde{U}(z,t), X_{ki}(z) \rangle, \qquad i = 1, 2.$$
 (4.19)

Applying the fractional operator to both sides of the first relation in (4.19), we have

$$\left({^{C}D_{0_{+},t}^{\alpha_{0}}} + \sum_{i=1}^{m} \mu_{i}{^{C}D_{0_{+},t}^{\alpha_{i}}}\right) \widetilde{U}_{ki}(t) = \left(\left({^{C}D_{0_{+},t}^{\alpha_{0}}} + \sum_{i=1}^{m} \mu_{i}{^{C}D_{0_{+},t}^{\alpha_{i}}}\right) \widetilde{U}(z,t), X_{ki}(z)\right),$$

By virtue of (1.5), we have the following equation

$${}^{C}D_{0_{+},t}^{\alpha_{0}}\widetilde{U}_{k,1}(t) + \sum_{i=1}^{m} \mu_{i}{}^{C}D_{0_{+},t}^{\alpha_{i}}\widetilde{U}_{k,1}(t) = -\lambda_{ki}\widetilde{U}_{ki}(t) + \widetilde{s}(t)h_{ki}(t),$$

The solution of the above equation is

$$\begin{split} \widetilde{U}_{ki}(t) &= \mathcal{E}_{\Xi,1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n})\widetilde{U}_{ki}(0) \\ &+ \sum_{m=1}^{n} \mu_{m} \mathcal{E}_{\Xi,\alpha_{0}-\alpha_{m}+1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}) \\ \widetilde{U}_{ki}(0) + \left\{ s(t)h_{ki}(t) \right\} * \mathcal{E}_{\Xi,\alpha_{0}+1}(t;\lambda_{ki},\mu_{1},\mu_{2},...,\mu_{n}). \end{split}$$

By equation (4.18), uniqueness of $\widetilde{s}(t)$ and completeness of $X_{ki}(z)$, we obtain

$$\widetilde{U}_{ki}(t) = 0, \quad \forall \ n \in \mathbb{N} \cup \{0\}.$$

and hence U(z,t) = V(z,t).

5. Conclusions

We considered two IPs for TFDE with involution along with one Dirichlet type and second boundary condition of non-local type. The eigenfunction expansion method is used to have series representation of the solutions of both IPs. The unique existence of the time dependent source term is obtained by using Banach fixed point theorem in IP-2. We checked existence and uniqueness results for both IPs by considering suitable over-specified and certain regularity conditions on given datum. In future, we consider an another interesting problem by considering multi-time Hilfer fractional derivatives in time with same spectral problem.

Acknowledgment: The author's express their sincere gratitude to Dr. Salman Amin Malik, for their valuable assistance in the manuscript. His insightful feedback and support significantly contributed to the quality of this work.

Conflicts of Interest: The authors declare that there are no conflicts of interest regarding the publication of this paper.

References

- [1] M.A. Sadybekov, G. Dildabek, M.B. Ivanova, On an Inverse Problem of Reconstructing a Heat Conduction Process from Nonlocal Data, Adv. Math. Phys. 2018 (2018), 8301656. https://doi.org/10.1155/2018/8301656.
- [2] A. Cabada, A.F. Tojo, Equations With Involutions, in: Workshop on Differential Equations 2014, Malá Morávka, Czech Republic, (2014).
- [3] M. Kirane, M.A. Sadybekov, A.A. Sarsenbi, On an Inverse Problem of Reconstructing a Subdiffusion Process From Nonlocal Data, Math. Methods Appl. Sci. 42 (2019), 2043–2052. https://doi.org/10.1002/mma.5498.
- [4] L. Salari, L. Rondoni, C. Giberti, R. Klages, A Simple Non-Chaotic Map Generating Subdiffusive, Diffusive, and Superdiffusive Dynamics, Chaos 25 (2015), 073113. https://doi.org/10.1063/1.4926621.
- [5] G. H. Weiss, Aspects and Applications of the Random Walk, North-Holland, New York, 1994.
- [6] B.D. Hughes, Random Walks and Random Environments, Clarendon Press, Oxford, 2002.

- [7] A. Mura, G. Pagnini, Characterizations and Simulations of a Class of Stochastic Processes to Model Anomalous Diffusion, J. Phys.: Math. Theor. 41 (2008), 285003. https://doi.org/10.1088/1751-8113/41/28/285003.
- [8] E. Nadal, E. Abisset-Chavanne, E. Cueto, F. Chinesta, On the Physical Interpretation of Fractional Diffusion, C. R. Mécanique 346 (2018), 581–589. https://doi.org/10.1016/j.crme.2018.04.004.
- [9] P. Paradisi, Fractional Calculus in Statistical Physics: The Case of Time Fractional Diffusion Equation, Commun. Appl. Ind. Math. 6 (2015), e-530.
- [10] R. Gorenflo, F. Mainardi, D. Moretti, P. Paradisi, Time Fractional Diffusion: a Discrete Random Walk Approach, Nonlinear Dyn. 29 (2002), 129–143. https://doi.org/10.1023/a:1016547232119.
- [11] E. Bakalis, F. Zerbetto, Time Fractional Diffusion Equations and Analytical Solvable Models, J. Phys.: Conf. Ser. 738 (2016), 012106. https://doi.org/10.1088/1742-6596/738/1/012106.
- [12] B. Baeumer, M. Kovács, M.M. Meerschaert, H. Sankaranarayanan, Reprint of: Boundary Conditions for Fractional Diffusion, J. Comput. Appl. Math. 339 (2018), 414–430. https://doi.org/10.1016/j.cam.2018.03.007.
- [13] E. Bakalis, F. Zerbetto, Time Fractional Diffusion Equations and Analytical Solvable Models, J. Phys.: Conf. Ser. 738 (2016), 012106. https://doi.org/10.1088/1742-6596/738/1/012106.
- [14] N. Shieh, On Time-Fractional Relativistic Diffusion Equations, J. Pseudo-Differential Oper. Appl. 3 (2012), 229–237. https://doi.org/10.1007/s11868-012-0049-6.
- [15] B. Ahmad, A. Alsaedi, M. Kirane, R.G. Tapdigoglu, An Inverse Problem for Space and Time Fractional Evolution Equations with an Involution Perturbation, Quaest. Math. 40 (2017), 151–160. https://doi.org/10.2989/16073606. 2017.1283370.
- [16] F.G. Waqar, S. Patel, C.M. Simon, A Tutorial on the Bayesian Statistical Approach to Inverse Problems, arXiv:2304.07610 (2023). https://doi.org/10.48550/ARXIV.2304.07610.
- [17] I.A. Kaliev, M.M. Sabitova, Problems of Determining the Temperature and Density of Heat Sources From the Initial and Final Temperatures, J. Appl. Ind. Math. 4 (2010), 332–339. https://doi.org/10.1134/s199047891003004x.
- [18] M. Kirane, S.A. Malik, M.A. Al-Gwaiz, An Inverse Source Problem for a Two Dimensional Time Fractional Diffusion Equation with Nonlocal Boundary Conditions, Math. Methods Appl. Sci. 36 (2012), 1056–1069. https://doi.org/10. 1002/mma.2661.
- [19] M. Kirane, S.A. Malik, Determination of an Unknown Source Term and the Temperature Distribution for the Linear Heat Equation Involving Fractional Derivative in Time, Appl. Math. Comput. 218 (2011), 163–170. https://doi.org/10.1016/j.amc.2011.05.084.
- [20] J. Wen, X. Ren, S. Wang, Simultaneous Determination of Source Term and the Initial Value in the Space-Fractional Diffusion Problem by a Novel Modified Quasi-Reversibility Regularization Method, Phys. Scr. 98 (2023), 025201. https://doi.org/10.1088/1402-4896/acaa68.
- [21] J. Wen, Z. Liu, S. Wang, Conjugate Gradient Method for Simultaneous Identification of the Source Term and Initial Data in a Time-Fractional Diffusion Equation, Appl. Math. Sci. Eng. 30 (2022), 324–338. https://doi.org/10.1080/27690911.2022.2075358.
- [22] J. Wen, Z. Liu, S. Wang, A Non-Stationary Iterative Tikhonov Regularization Method for Simultaneous Inversion in a Time-Fractional Diffusion Equation, J. Comput. Appl. Math. 426 (2023), 115094. https://doi.org/10.1016/j.cam. 2023.115094.
- [23] J. Wen, Y. Wang, The Quasi-Reversibility Regularization Method for Backward Problems of the Time-Fractional Diffusion-Wave Equation, Phys. Scr. 98 (2023), 095250. https://doi.org/10.1088/1402-4896/acf0f8.
- [24] T. Wei, K. Liao, Identifying a Time-Dependent Zeroth-Order Coefficient in a Time-Fractional Diffusion-Wave Equation by Using the Measured Data at a Boundary Point, Appl. Anal. 101 (2021), 6522–6547. https://doi.org/10. 1080/00036811.2021.1932834.
- [25] T. Wei, Y. Zhang, D. Gao, Identification of the Zeroth-Order Coefficient and Fractional Order in a Time-fractional Reaction-diffusion-wave Equation, Math. Methods Appl. Sci. 46 (2022), 142–166. https://doi.org/10.1002/mma.8499.

- [26] Z. Ruan, Q. Hu, W. Zhang, Identification of a Time-Dependent Control Parameter for a Stochastic Diffusion Equation, Comput. Appl. Math. 40 (2021), 201. https://doi.org/10.1007/s40314-021-01598-0.
- [27] Z. Ruan, Z. Wang, A Backward Problem for Distributed Order Diffusion Equation: Uniqueness and Numerical Solution, Inverse Probl. Sci. Eng. 29 (2020), 418–439. https://doi.org/10.1080/17415977.2020.1795152.
- [28] A. Ashyralyev, A. Sarsenbi, Well-posedness of a Parabolic Equation with Nonlocal Boundary Condition, Bound. Value Probl. 2015 (2015), 38. https://doi.org/10.1186/s13661-015-0297-5.
- [29] M. Kirane, N. Al-Salti, Inverse Problems for a Nonlocal Wave Equation with an Involution Perturbation, J. Nonlinear Sci. Appl. 09 (2016), 1243–1251. https://doi.org/10.22436/jnsa.009.03.49.
- [30] A. Ashyralyev, A. Sarsenbi, Well-posedness of a Parabolic Equation with Involution, Numer. Funct. Anal. Optim. 38 (2017), 1295–1304. https://doi.org/10.1080/01630563.2017.1316997.
- [31] N.H. Tuan, M. Kirane, L.V.C. Hoan, L.D. Long, Identification and Regularization for Unknown Source for a Time-Fractional Diffusion Equation, Comput. Math. Appl. 73 (2017), 931–950. https://doi.org/10.1016/j.camwa.2016.10.002.
- [32] I. Orazov, M.A. Sadybekov, On a Class of Problems of Determining the Temperature and Density of Heat Sources Given Initial and Final Temperature, Sib. Math. J. 53 (2012), 146–151. https://doi.org/10.1134/s0037446612010120.
- [33] N.I. Ivanchov, Inverse Problems for the Heat-Conduction Equation with Nonlocal Boundary Conditions, Ukr. Math. J. 45 (1993), 1186–1192. https://doi.org/10.1007/bf01070965.
- [34] I.A. Kaliev, M.F. Mugafarov, O.V. Fattahova, Inverse Problem for Forward-Backward Parabolic Equation With Generalized Conjugation Conditions, Ufa Math. J. 3 (2011), 33–41.
- [35] F. Kanca, M.I. Ismailov, The Inverse Problem of Finding the Time-Dependent Diffusion Coefficient of the Heat Equation From Integral Overdetermination Data, Inverse Probl. Sci. Eng. 20 (2011), 463–476. https://doi.org/10. 1080/17415977.2011.629093.
- [36] A.B. Kostin, Counterexamples in Inverse Problems for Parabolic, Elliptic, and Hyperbolic Equations, Comput. Math. Phys. 54 (2014), 797–810. https://doi.org/10.1134/s0965542514020092.
- [37] F. Kanca, Inverse Coefficient Problem of the Parabolic Equation with Periodic Boundary and Integral Overdetermination Conditions, Abstr. Appl. Anal. 2013 (2013), 659804. https://doi.org/10.1155/2013/659804.
- [38] D. Lesnic, S.A. Yousefi, M. Ivanchov, Determination of a Time-Dependent Diffusivity From Nonlocal Conditions, J. Appl. Math. Comput. 41 (2012), 301–320. https://doi.org/10.1007/s12190-012-0606-4.
- [39] L. Miller, M. Yamamoto, Coefficient Inverse Problem for a Fractional Diffusion Equation, Inverse Probl. 29 (2013), 075013. https://doi.org/10.1088/0266-5611/29/7/075013.
- [40] G. Li, D. Zhang, X. Jia, M. Yamamoto, Simultaneous Inversion for the Space-Dependent Diffusion Coefficient and the Fractional Order in the Time-Fractional Diffusion Equation, Inverse Probl. 29 (2013), 065014. https://doi.org/10. 1088/0266-5611/29/6/065014.
- [41] N.H. Tuan, D.N.D. Hai, L.D. Long, V.T. Nguyen, M. Kirane, On a Riesz–feller Space Fractional Backward Diffusion Problem with a Nonlinear Source, J. Comput. Appl. Math. 312 (2017), 103–126. https://doi.org/10.1016/j.cam.2016. 01.003.
- [42] , Existence Results on Mittag-Leffler Kernel of Fractional Integro-Differential Inclusion Problem Under Boundary Conditions, Prog. Fract. Differ. Appl. 10 (2024), 149–160. https://doi.org/10.18576/pfda/100114.
- [43] A. Erdogan, D. Kusmangazinova, I. Orazov, M. Sadybekov, On One Problem for Restoring the Density of Sources of the Fractional Heat Conductivity Process with Respect to Initial and Final Temperatures, Bull. Karaganda Univ. Math. 91 (2018), 31–44. https://doi.org/10.31489/2018m3/31-44.
- [44] M. Kirane, B. Samet, B.T. Torebek, Determination of an Unknown Source Term Temperature Distribution for the Sub-Diffusion Equation at the Initial and Final Data, Electron. J. Differ. Equ. 2017 (2017), 257.
- [45] O. Taki-Eddine, B. Abdelfatah, On Determining the Coefficient in a Parabolic Equation With Nonlocal Boundary and Integral Condition, Electron. J. Math. Anal. Appl. 6 (2018), 94–102.

- [46] M. Ali, S. Aziz, S.A. Malik, Inverse Problem for a Space-time Fractional Diffusion Equation: Application of Fractional Sturm-liouville Operator, Math. Methods Appl. Sci. 41 (2018), 2733–2747. https://doi.org/10.1002/mma. 4776.
- [47] S.G. Samko, A.A. Kilbas, O.I. Marichev, Fractional Integrals and Derivatives: Theory and Applications, Gordon and Breach, 1993.
- [48] Y. Luchko, R. Gorenflo, An Operational Method for Solving Fractional Differential Equations With the Caputo Derivatives, Acta Math. Vietnam 24 (1999), 207–233.
- [49] Z. Li, Y. Liu, M. Yamamoto, Initial-boundary Value Problems for Multi-Term Time-Fractional Diffusion Equations with Positive Constant Coefficients, Appl. Math. Comput. 257 (2015), 381–397. https://doi.org/10.1016/j.amc.2014. 11.073.
- [50] S.A. Malik, A. Ilyas, A. Samreen, Simultaneous Determination of a Source Term and Diffusion Concentration for a Multi-Term Space-Time Fractional Diffusion Equation, Math. Model. Anal. 26 (2021), 411–431. https://doi.org/10. 3846/mma.2021.11911.
- [51] M. Ali, S. Aziz, S.A. Malik, Inverse Source Problems for a Space–time Fractional Differential Equation, Inverse Probl. Sci. Eng. 28 (2019), 47–68. https://doi.org/10.1080/17415977.2019.1597079.